A multi-class vehicular flow model for aggressive drivers.

R. M. Velasco Aggressive multi-class model

A multi-class vehicular flow model for aggressive drivers.

W. Marques Jr., R. M. Velasco, A. R. Méndez

27 October, 2015

Universidade Federal do Paraná, Brazil. Universidad Autónoma Metropolitana-Iztapalapa, México. Universidad Autónoma Metropolitana-Cuajimalpa, México. **Abstract.** The kinetic theory approaches to vehicular traffic modeling have given very good results in the understanding of the dynamical phenomena involved. In this work we deal with the kinetic approach modeling of a traffic situation where there are many classes of aggressive drivers. Their aggressivity is characterized through their relaxation times. The reduced Paveri-Fontana equation is taken as a starting point to set the model. The kinetic equation is taken to write the balance equations for the density

and the average speed in each drivers class. In this model we consider that each class of drivers preserve the corresponding aggressivity, in such a way that there will be no adaptation effects. It means that the number of drivers in a class is conserved. Some characteristics of the model are explored with the usual methods.

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Introduction.

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$$\frac{\partial f_i(c)}{\partial t} + c \frac{\partial f_i(c)}{\partial x} + \frac{\partial}{\partial c} \left(\frac{c_0(c) - c}{\tau_i} f_i(c) \right) = q \sum_j \left[\rho_i f_j(c) \xi_i(c) + \rho_j f_i(c) \psi_j(c) \right],$$
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The interaction is separated according to active $\psi_i(c)$ or passive $\xi_i(c)$ interaction terms.

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$$\psi_i(c) = \int_{w < c} (w - c) \frac{f_i(w)}{\rho_i} dw, \quad i = 1, 2 \quad \text{active}, \tag{2}$$

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Both give a contribution to the balance in the kinetic equation [7].

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The macroscopic variables

We define the densities and average speeds as

$$\rho_i(x,t) = \int f_i(c) dc, \qquad v_i(x,t) = \int c \frac{f_i(c)}{\rho_i} dc = \langle c \rangle_i, \qquad (4)$$

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where $\langle ... \rangle_i$ means the average over the f_i distribution function. It should be noticed that $\langle \xi_j \rangle_i - \langle \psi_i \rangle_j = v_j - v_i$ is satisfied.

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$$\frac{\partial \rho_i v_i}{\partial t} + \frac{\partial}{\partial x} (\rho_i v_i^2 + \mathcal{P}_i) = \frac{\rho_i}{\tau_i} (V_i^0 - v_i) + q \sum_j \rho_i \rho_j [\langle c\xi_i \rangle_j + \langle c\psi_j \rangle_i], \quad (6)$$

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$$\frac{\partial \rho_i \mathbf{v}_i}{\partial t} + \frac{\partial}{\partial x} (\rho_i \mathbf{v}_i^2 + \mathcal{P}_i) = \frac{\rho_i}{\tau_i} (\mathbf{V}_i^0 - \mathbf{v}_i) + q \sum_j \rho_i \rho_j [\langle c\xi_i \rangle_j + \langle c\psi_j \rangle_i], \quad (6)$$

where $V_i^0(x, t)$ comes from the average of the desired speed now taken over the instantaneous speed c, and $\mathcal{P}_i = \int (c - v_i)^2 f_i(c) dc = \rho_i \Theta_i$ is the *i*-class traffic pressure and Θ_i the speed variance. First of all we consider aggressive drivers characterized by the parameter ω in such a way that the average desired speed is assumed as $V_i^0(c) = \omega v_i$, $\omega > 1$.

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It should be noticed that $\rho_i(x, t)$, $v_i(x, t)$ in Eq. (7) are local variables.

Now the interaction terms

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$$\mathcal{I}_{ij}(c) = \langle c\xi_i \rangle_j + \langle c\psi_j \rangle_i, \tag{8}$$

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The active and passive interactions between different classes for $(i \neq j)$ can be written in a closed form, however we have found that a simpler and very good approximation for them

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$$\mathcal{I}_{ij} = v_i v_j \left\{ 1 - \left(\frac{\alpha + 1}{\alpha} \frac{v_j}{v_i} \right) \mathcal{H}(v_i - v_j) + \left(1 - \frac{\alpha + 1}{\alpha} \frac{v_i}{v_j} \right) \mathcal{H}(v_i - v_j) \right\},$$
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$$\mathcal{I}_{11} = -\Theta_1, \ \mathcal{I}_{22} = -\Theta_2, \ \mathcal{I}_{12} = 0, \ \mathcal{I}_{21} \simeq -(v_2 - v_1)^2.$$
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$$\Theta_i(x,t) = \frac{\mathsf{v}_i^2}{\alpha},\tag{12}$$

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where μ_i is a constant analogous of a viscosity and it is different for each class of drivers. It should be said that this hypothesis can be justified in terms of the kinetic model [6, 8].

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$$\frac{\partial \rho_1}{\partial t} + \frac{\partial \rho_1 v_1}{\partial x} = 0, \tag{14}$$

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$$\frac{\partial v_1}{\partial t} + v_1 \frac{\partial v_1}{\partial x} = -\frac{1}{\rho_1} \frac{\partial \mathcal{P}_1}{\partial x} + \frac{v_1^* - v_1}{\tau_1},$$
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where

$$v_1^* = wv_1 + \tau_1(1-p)\rho_1 \mathcal{I}_{11} + \tau_1(1-p)\rho_2 \mathcal{I}_{12}, \qquad (18)$$

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$$\frac{\partial v_2}{\partial t} + v_2 \frac{\partial v_2}{\partial x} = -\frac{1}{\rho_2} \frac{\partial \mathcal{P}_2}{\partial x} + \frac{v_2^* - v_2}{\tau_2}$$
(17)

where

$$v_1^* = wv_1 + \tau_1(1-p)\rho_1 \mathcal{I}_{11} + \tau_1(1-p)\rho_2 \mathcal{I}_{12}, \qquad (18)$$

$$v_{2}^{*} = wv_{1} + \tau_{2}(1-p)\rho_{1}\mathcal{I}_{21} + \tau_{2}(1-p)\rho_{2}\mathcal{I}_{22}, \qquad (19)$$

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(21)

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where the perturbation has been expanded in modes with a wave vector k and a complex frequency called σ in such a way that the stability condition for the equilibrium state is given by $\mathcal{R}e\sigma > 0$.

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$$\begin{pmatrix} -\sigma + ikv_1^e & 0 & ik\rho_1^e & 0 \\ \frac{(v_1^e)^2 ik}{\alpha\rho_1^e} - \frac{\gamma_{11}^e}{\tau_1} & 0 & -\sigma + \frac{\alpha+2}{\alpha}ikv_1^e + \frac{\mu_1}{\rho_1^e}k^2 + \frac{1}{\tau_1} & 0 \\ 0 & -\sigma + ikv_2^e & 0 & ik\rho_2^e \\ -\frac{a_1}{\tau_2} & \frac{ik(v_2^e)^2}{\alpha\rho_2^e} - \frac{a_2}{\tau_2} & -\frac{a_3}{\tau_2} & -\sigma + \frac{\alpha+2}{\alpha}ikv_2^e + \frac{\mu_2k^2}{\rho_2^e} + \frac{1-a_4}{\tau_2} \end{pmatrix},$$
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where its determinant must vanish.

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$$\begin{pmatrix} -\sigma + ikv_1^e & 0 & ik\rho_1^e & 0 \\ \frac{(v_1^e)^{2ik}}{\alpha\rho_1^e} - \frac{\gamma_{11}^e}{\tau_1} & 0 & -\sigma + \frac{\alpha+2}{\alpha}ikv_1^e + \frac{\mu_1}{\rho_1^e}k^2 + \frac{1}{\tau_1} & 0 \\ 0 & -\sigma + ikv_2^e & 0 & ik\rho_2^e \\ -\frac{\vartheta_1}{\tau_2} & \frac{ik(v_2^e)^2}{\alpha\rho_2^e} - \frac{\vartheta_2}{\tau_2} & -\frac{\vartheta_3}{\tau_2} & -\sigma + \frac{\alpha+2}{\alpha}ikv_2^e + \frac{\mu_2k^2}{\rho_2^e} + \frac{1-\vartheta_4}{\tau_2} \end{pmatrix},$$
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Due to the fact that the macroscopic equations are valid in a kind of hydrodynamical limit $(k \rightarrow 0)$, we will expand the roots in the dispersion relation around k = 0 and take terms up to order k^2 ,

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$$\sigma = \sigma_0 + k\sigma_1 + k^2\sigma_2 + \mathcal{O}(k^3), \qquad (24)$$

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and there will be four different roots, which will be called as Σ_i .

$$\Sigma_{1} = ikc_{1} + \frac{k^{2}\tau_{1}}{\alpha\rho_{1}^{e}} [c_{1}^{2} - (\alpha + 1)(\gamma_{11}^{e}\rho_{1}^{e})^{2}]$$
(25)

$$\Sigma_{2} = \frac{1}{\tau_{1}} + \frac{ik}{\alpha} (2v_{1}^{e} + \alpha v_{1}^{e} - \alpha \gamma_{11}^{e} \rho_{1}^{e}) + \frac{k^{2}}{\alpha \rho_{1}^{e}} \Big[\alpha \mu_{1} - \tau_{1} \rho_{1}^{e} ((v_{1}^{e})^{2} + 2\rho_{1}^{e} v_{1}^{e} \gamma_{11}^{e} - \alpha (\rho_{1}^{e} \gamma_{11}^{e})^{2}) \Big] + \dots$$
(26)

$$\Sigma_{3} = \frac{ik}{a_{4} - 1} \left[(a_{4} - 1)v_{2}^{e} - a_{2}\rho_{2}^{e} \right] + k^{2}\Sigma_{32}$$

$$\Sigma_{32} = -\frac{\tau_{2}}{\alpha(a_{4} - 1)^{3}} \left[(a_{4} - 1)^{2}(v_{2}^{e})^{2} - a_{2}\rho_{2}^{e} \left(2v_{2}^{e}(a_{4} - 1) + \alpha a_{2}\rho_{2}^{e} \right) \right], \qquad (27)$$

$$\Sigma_{4} = \frac{1 - a_{4}}{\tau_{2}} + \frac{ik}{\alpha(a_{4} - 1)} \Big[(a_{4} - 1)v_{2}^{e}(2 + \alpha) + \alpha a_{2}\rho_{2}^{e} \Big] + k^{2}\Sigma_{42}$$

$$\Sigma_{42} = \frac{1}{(a_{4} - 1)^{3}\alpha\rho_{2}^{e}} \Big[(a_{4} - 1)\rho_{2}^{e}v_{2}^{e}\tau_{2} \Big((a_{4} - 1)v_{2}^{e} - 2a_{2}\rho_{2}^{e} \Big) + \alpha \Big[(a_{4} - 1)^{3}\mu_{2} - \rho_{2}^{e}\tau_{2} (a_{2}\rho_{2}^{e})^{2} \Big) \Big]$$
(28)

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In this case the fundamental diagram chosen for v_1^e determines its behavior and the corresponding stability property. If we take the Greenshields diagram we obtain a stability region for $\rho_1^e < 43 \text{ veh/km}$, which means that we are at low density.

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Lastly, the mode corresponding to the class 2 density is also a propagating mode, its speed being

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And the real part in Σ_{32} must be positive to obtain stability in the mode corresponding to the density in class 2.



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Casa abierta al tiempo

R. M. Velasco Aggressive multi-class model

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