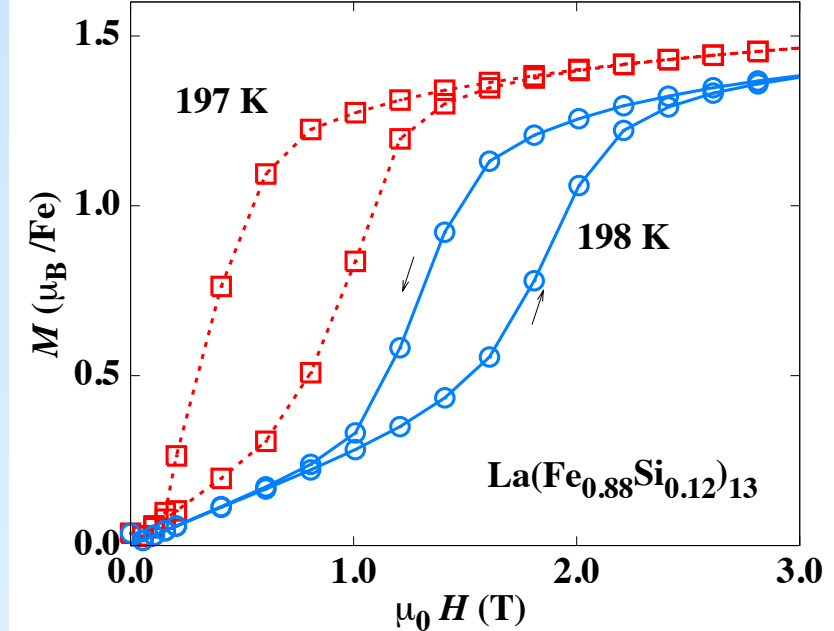
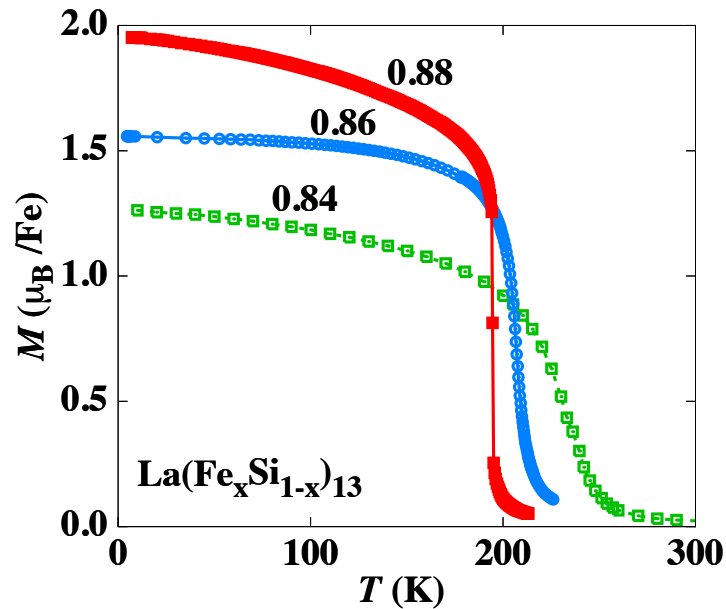


Influence of Magnetic Field on Nucleation of Thermally-induced Phase Transition in $\text{La}(\text{Fe}_{0.88}\text{Si}_{0.12})_{13}$

A. Fujita and M. Kano
Department of Materials Science,
Tohoku University, Japan

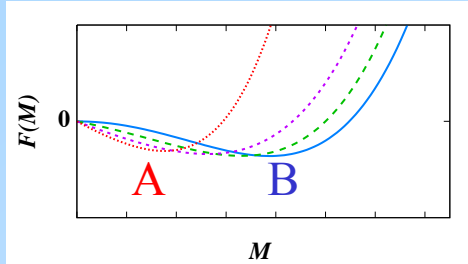


A. Fujita, Y. Akamatsu and K. Fukamichi, J. Appl. Phys. **85** (1999) 4756.

Why kinetics ?

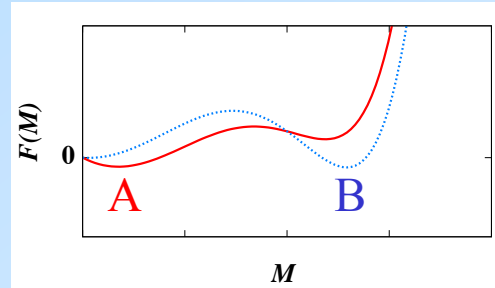
For the 1st order phase transition

* No homogeneous ΔS_m and ΔT_{ad}



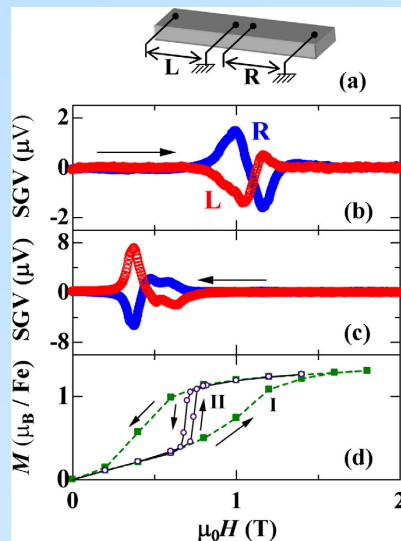
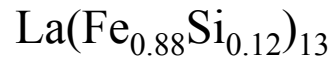
* 2nd-order

Route from state A to state B is continuous. "Isothermal" or "Adiabatic" path : OK

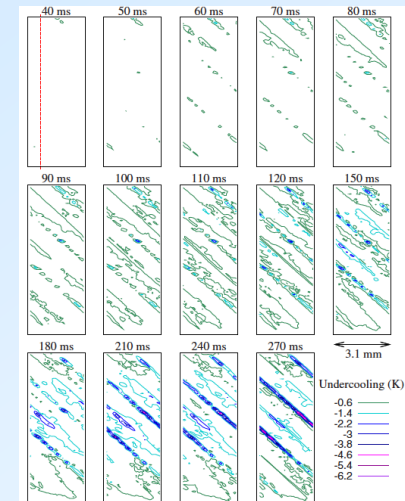
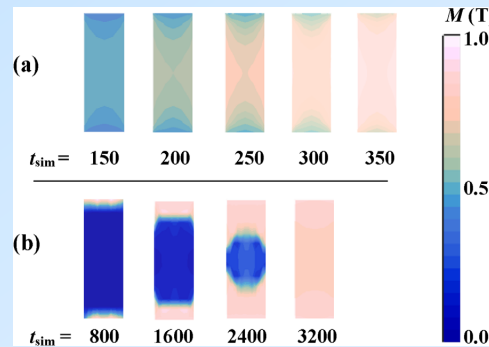
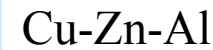


* 1st-order

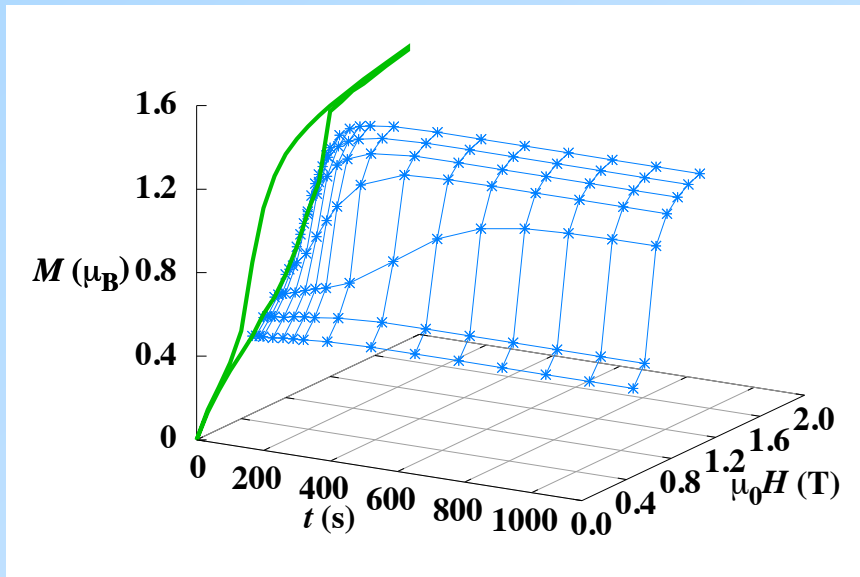
Even if A and B have same S or T inter-state transition is terminated by energy barrier



Fujita et al., *Apl. Phys. Lett.* **102** (2013) 041913.

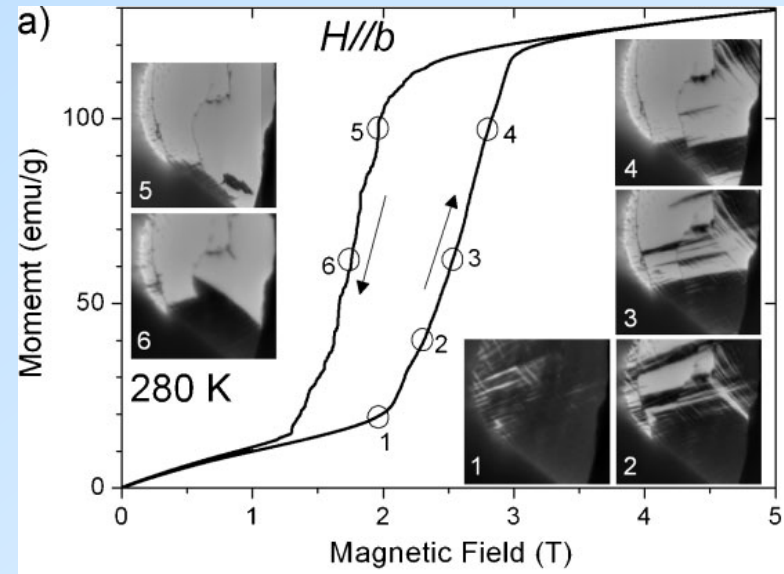


Vives, Mañosa, Planes et al.
Apl. Phys. Lett. **98** (2011) 011902.



$\text{La}(\text{Fe}_{0.88}\text{Si}_{0.12})_{13}$
 Yako and Fujita, IEEE Trans. Magn.,
 47(2011) 2482 ,

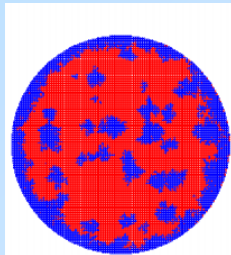
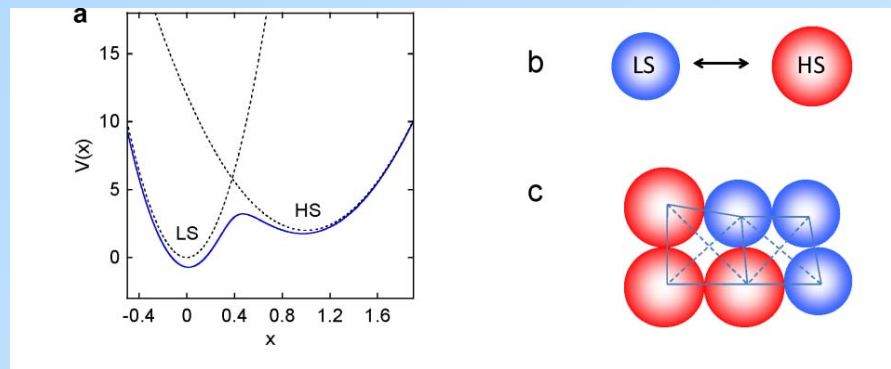
“Iso-field”



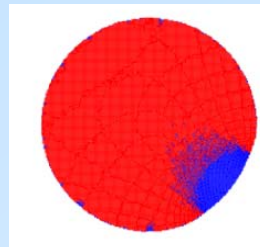
$\text{Gd}_5\text{Ge}_2\text{Si}_2$
 Moor et al, Adv. Mater., 21(2009), 3780

“ a-field ?? ”

Spin crossover



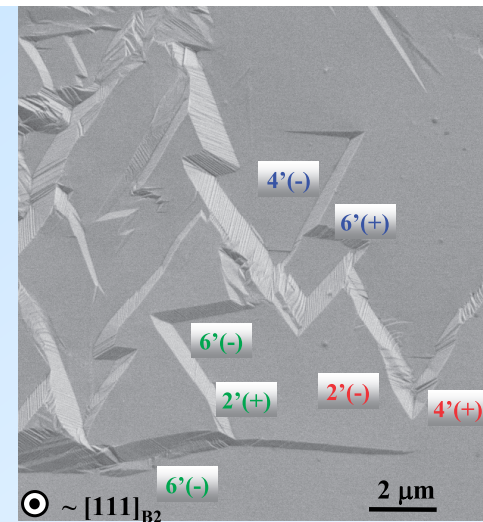
SRO model



SRO+LRO

Nishino, Miyashita et al
Sci. Rep. 1 (2011) 162

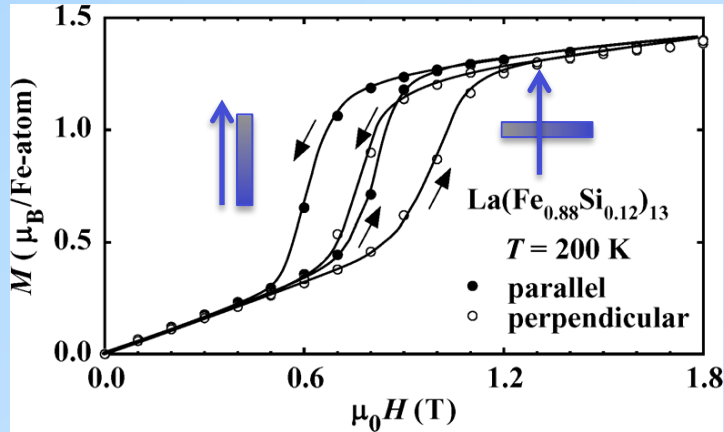
Long-distance strain-connected morphology in Ti-Ni martensite



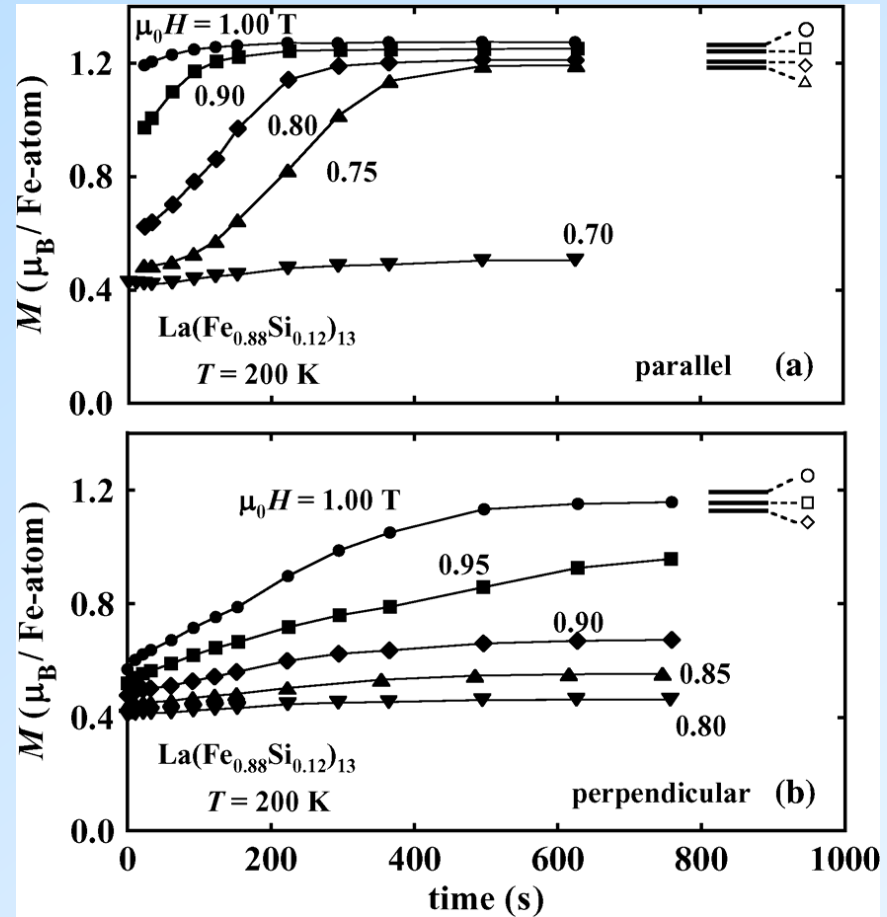
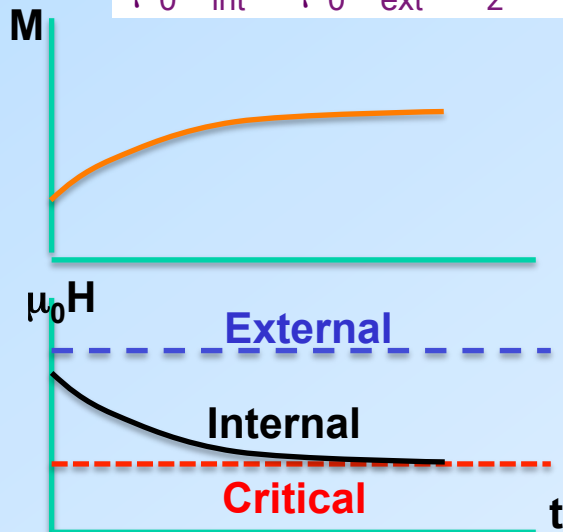
Nishida et al, Phil. Mag. 92 (2012) 221

“Shape anisotropy” in kinetics

1mm * 1mm * 4mm
rectangular parallelepiped



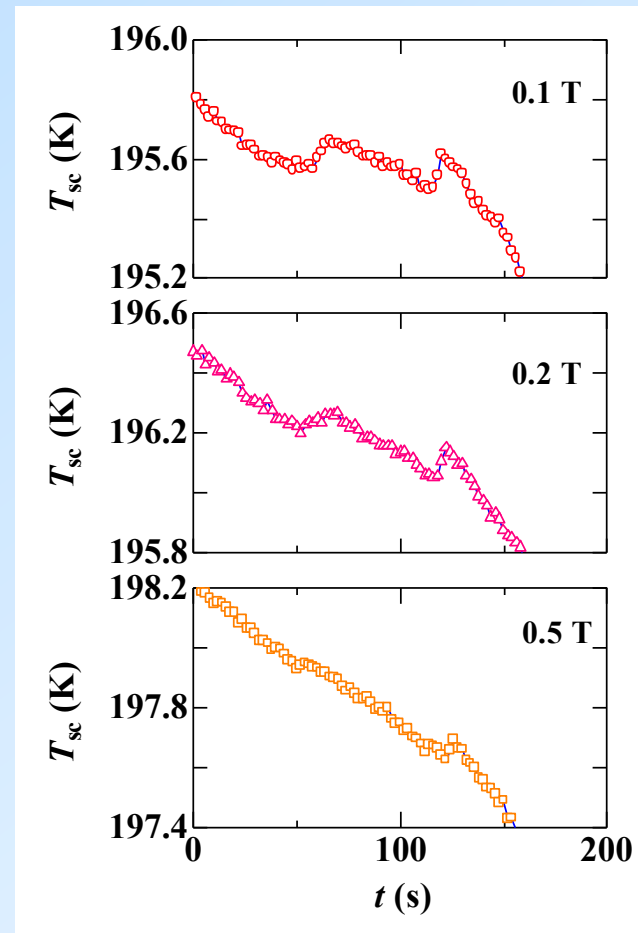
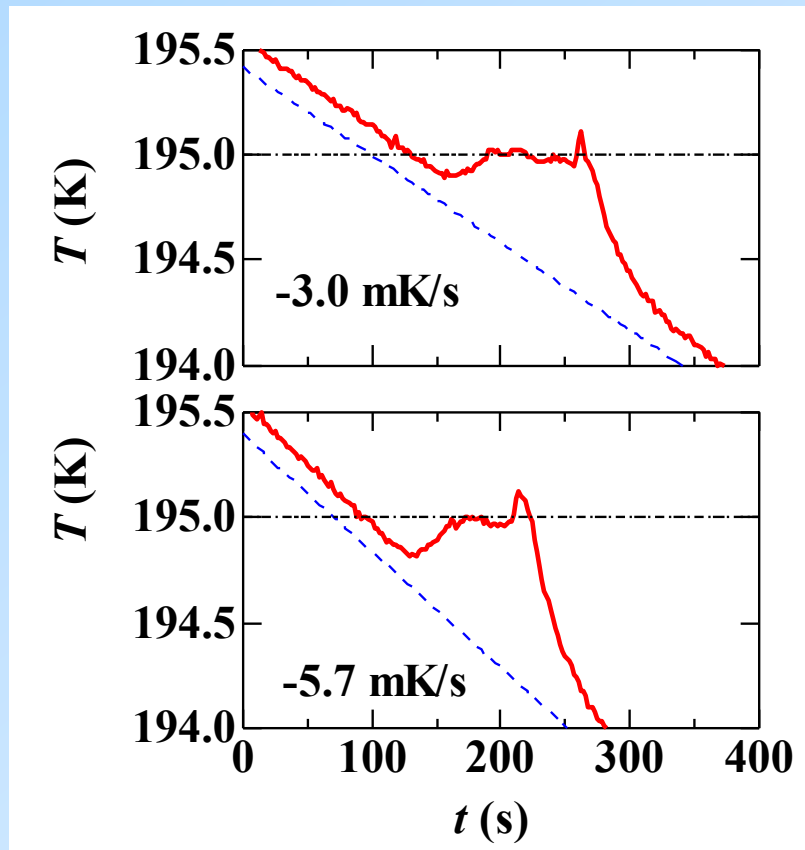
$$\mu_0 H_{\text{int}} = \mu_0 H_{\text{ext}} - D_z M$$



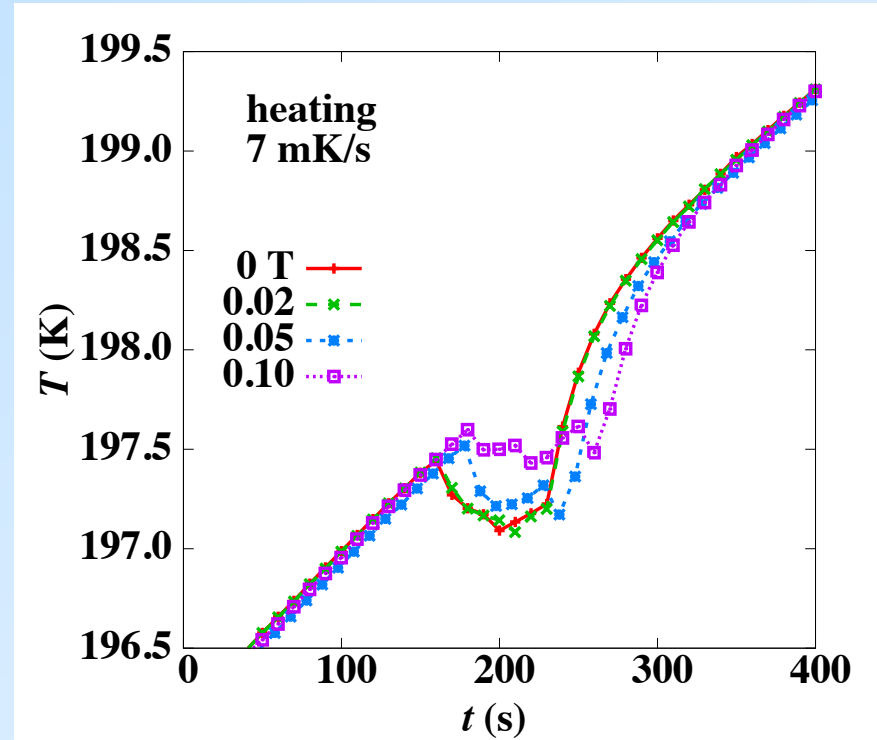
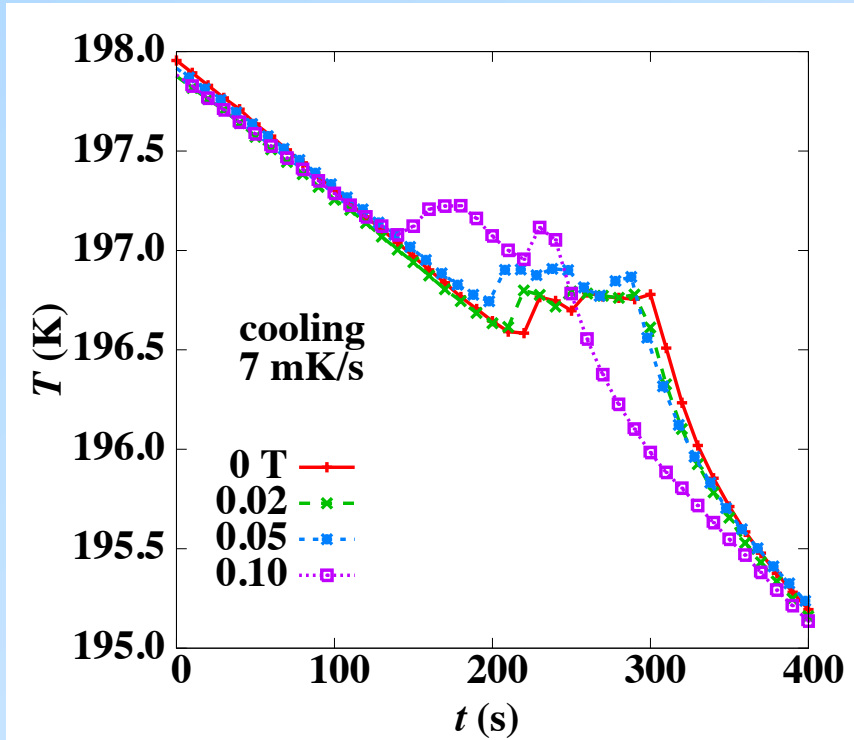
Yako and Fujita, IEEE Trans. Magn. 47 (2011) 2482

Value of unfinished M corresponds to $\mu_0 H_C - \mu_0 H_{\text{int}}$ by using “bulk” D_z value (instead of nuclei-shape one)

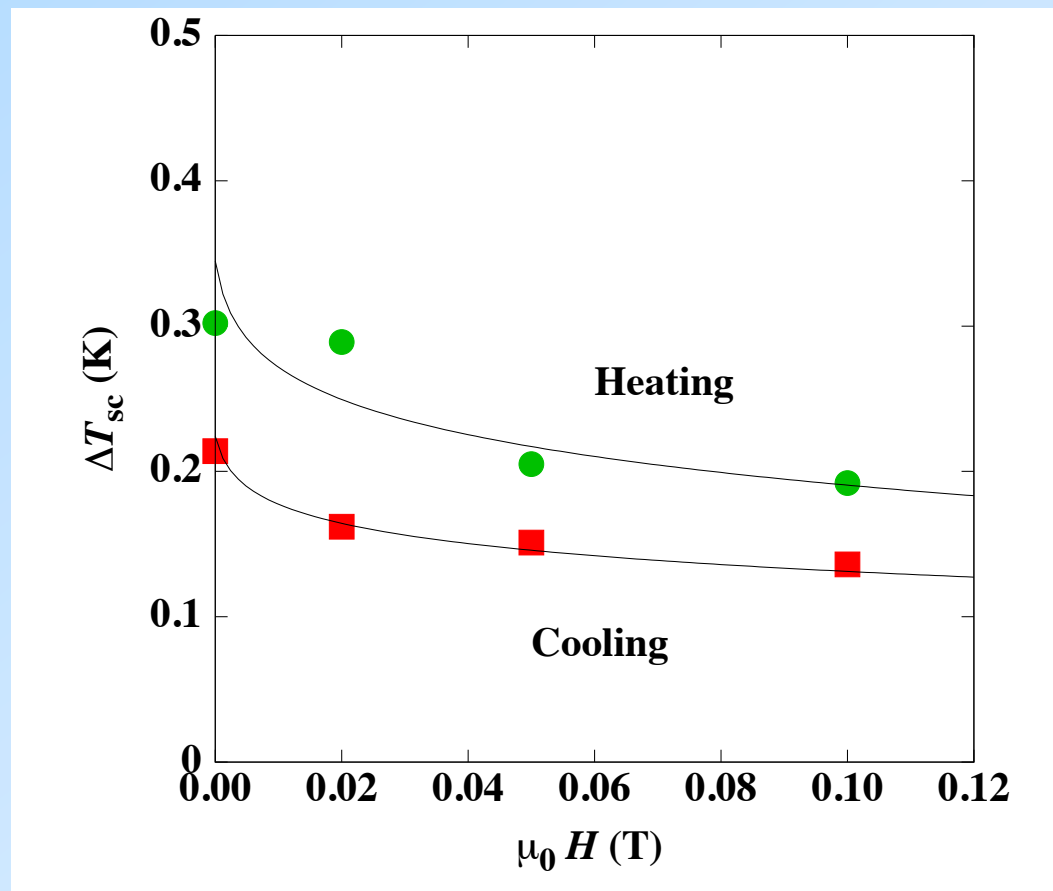
Supercooling feature



Supercooling(heating) profile under mag. field



Magnitude of undercooling



Classical model

$$\frac{2r_c}{\Delta T_{sc}} = A \exp\left(-\frac{\Delta G}{k_B T}\right)$$

$$\Delta G = \frac{16\pi\gamma^3}{3L^2} \left(\frac{T_C}{\Delta T_{sc}}\right)^2$$

L : latent heat γ : surface energy
 T_C : Curie temp. k_B : Boltzmann const.
 ΔT_{sc} : Magnitude of under cooling
 r_c : Temperature sweeping speed
 A : pre-exponential factor

Large difference of dipole distribution ???

Para. → Ferro.

Nucleation first occurs, then *magnetic domain* is formed after enough growth of the Ferro. region

Ferro. → Para.

Nucleation occurs in the Ferro. region including *magnetic domain*.

Magnetic fluctuations ???

Incubation from embryo to nuclei is reflects magnetic fluctuations

Magnetic field suppress the paramagnetic fluctuations

Conclusion

Nucleation behavior is characterized under magnetic field in $\text{La}(\text{Fe}_{0.88}\text{Si}_{0.12})_{13}$

1. Supercooling (heating) behavior shows different characteristics in cooling and heating processes.
2. By applying external magnetic field, magnitude of undercooling becomes smaller in both the cooling and heating process.
3. These change caused by magnetic field is explained by change in surface energy loss of nuclei in the conventional model.

Magnetostatic information, such as flux distribution of nuclei surrounded by magnetic domains, is necessary to reveal the nucleation-growth properties in magnetic fields (the AMR situations).