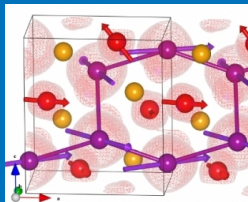
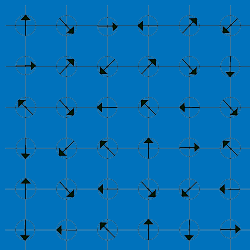


Disordered Local Moments, the Magnetocaloric Effect and Metamagnetism



Julie Staunton
University of Warwick

Acknowledgements

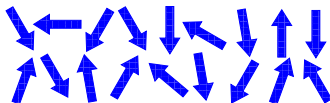
Manuel dos Santos Dias, Rudra Banerjee, Jonathan Peace (Warwick)
Zsolt Gercsi and Karl Sandeman (Imperial College)



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- For magnets these are **local moments**: $\{\hat{e}_i\}$



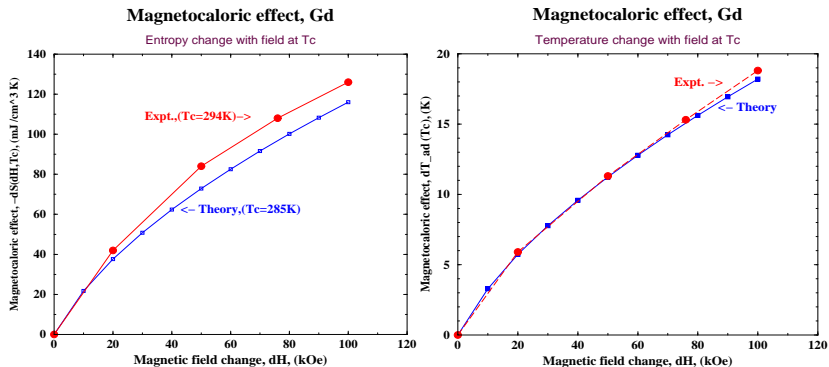
- $P(\{\hat{e}_i\}) = \frac{\exp[-\beta\Omega(\{\hat{e}_i\})]}{\prod_j \int d\hat{e}_j \exp[-\beta\Omega(\{\hat{e}_i\})]}$ where $\Omega(\{\hat{e}_i\})$ is the electronic grand potential from SDFT.

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- 'First-Principles' Mean Field Theory (DLM picture), averaging using techniques adapted from theory for electrons in disordered systems (CPA).

$$P_k(\hat{e}_k) = \frac{\exp[\beta\vec{h}_k \cdot \hat{e}_k]}{\int d\hat{e}_k \exp[-\beta\vec{h}_k \cdot \hat{e}_k]}, \quad \vec{m}_k = \int \hat{e}_k P_k(\hat{e}_k) d\hat{e}_k .$$

Ab-initio modelling of the MagnetoCaloric Effect in Gadolinium



Experimental results from K.Gschneidner Jr. et al., Rep.Prog.Phys. 68, (2005), 1479-1539

Magnetic interactions and structure from the DLM

- Free energy function

$$F(\{\vec{m}_i\}, \vec{B}, T) = \bar{E}(\{\vec{m}_i\}, \vec{B}) - T[\bar{S}_{mag}(\{\vec{m}_i\}, \vec{B}) + \bar{S}_{elec}(\{\hat{m}_i\}, \vec{B}, T) - \vec{B} \cdot \sum_i \mu_i \vec{m}_i]$$

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- The magnetic entropy

$$\bar{S}_{mag} = -k_B \sum_i \int P_i(\hat{e}_i) \ln P_i(\hat{e}_i) d\hat{e}_i$$

and electronic entropy

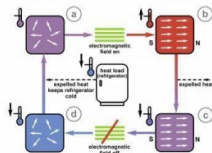
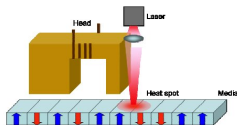
$$\bar{S}_{elec} \approx \frac{\pi^2 k_B T}{6} n(E_f; \{\vec{m}_i, \vec{B}\}),$$

hence MCE.

- Find Free energy by finding lowest $F(\{\vec{m}_i\}, \vec{B}, T)$ for several (FM, AFM, ...) magnetic states.

Metamagnetism

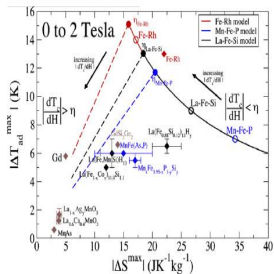
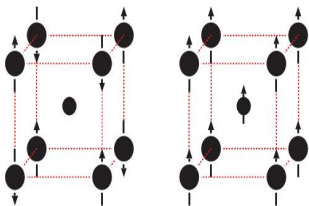
- Field or temperature induced transition from antiferromagnetic to ferromagnetic state.
- Second order, first order phase transition. Tricriticality, T_t, B_t .
- Heat assisted magnetic recording. Magnetic refrigeration.



Work near T_t, B_t .

- Material with competing ferromagnetic and anti-ferromagnetic interactions. Often linked with magnetostructural effect.
- Need to control for accessible and useful T_t, B_t .

Fe-Rh - a 'two-faced' magnetic alloy



On the left, the magnetic states of the ordered B2 (CsCl) alloy **Fe-Rh**.

On the right, experimental data for $|\Delta T_{ad}^{max}|$ vs. $|\Delta S^{max}|$ for several room temperature magnetic refrigerants, taken from viewpoint paper by K. G. Sandeman (*Scripta Mat.* **67**, 566-571, (2012)).

FeRh: A little compositional disorder goes a long way . . .

- The $\text{Fe}_{50}\text{Rh}_{50}$ solid solution orders into a B2 alloy at $T \approx 1600\text{K}$. Above $T = 0\text{K}$, the composition is $\text{Fe}_{(100-x)}\text{Rh}_x\text{-Rh}_{(100-x)}\text{Fe}_x$, where $x \neq 0$, the ordering incomplete.
- Away from stoichiometry and where compositional ordering is not complete, there can be Fe atoms on 'Rh' sites.
- Dramatic effect on magnetic properties, phase coexistence and broadening of 1st order transition. For $\text{Fe}_{49}\text{Rh}_{51}$ expt. finds a FM-AF transition at 370K ($T_c = 670\text{K}$) with $|\Delta S^{\max}| = 22.5 \text{ J K}^{-1} \text{ Kg}^{-1}$ at 2T.

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- DLM Theory calculates a FM-AF transition for perfect B2 Fe-Rh at $T_t = 495\text{K}$ ($T_c = 773\text{K}$). For 2T it finds $|\Delta S^{\max}| = 22.2 \text{ J K}^{-1} \text{ Kg}^{-1}$. **40 % of this is from electronic entropy.**

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- Incomplete B2 order: Swapping just **2% of Fe with Rh** causes T_t to drop to 208K ($T_c = 859\text{K}$). At **4%** FM-AF transition has vanished.
- Off stoichiometry: For $\text{Fe}_{96}\text{Rh}_4\text{-Rh}$, $T_t = 549\text{K}$ ($T_c = 700\text{K}$), no transition for $\text{Fe-Rh}_{96}\text{Fe}_4$ ($T_c = 1008\text{K}$).
- For $\text{Fe}_{97}\text{Rh}_3\text{-Rh}_{99}\text{Fe}_1$ ($\text{Fe}_{49}\text{Rh}_{51}$) $T_t = 415\text{K}$ ($T_c = 815\text{K}$). $|\Delta S^{\max}| = 20.7 \text{ J K}^{-1} \text{ Kg}^{-1}$ at 2T.

- **Local spin model** of helical AFM in external magnetic field (pairwise exchange interactions and mean field theory):
 - **No magnetic anisotropy** - spins cant smoothly into a conical spiral towards the field's direction, **2nd. order phase transition.**
 - **With magnetic anisotropy** - **first order transition** into fan structure, spins oscillating about field direction or FM state.

Metamagnetism and Helical Antiferromagnetic CoMnSi

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- **Itinerant electron spin model**. Landau free energy expansion, coefficients describe Stoner and spin fluctuation collective electron effects.
 - Free energy difference $\Delta F \approx$

$$a(\vec{m}_0, \vec{B}, T)|\Delta m_Q|^2 + b(\vec{m}_0, \vec{B}, T)|\Delta m_Q|^4 + c(\vec{m}_0, \vec{B}, T)|\Delta m_Q|^6$$

- Tricritical point:

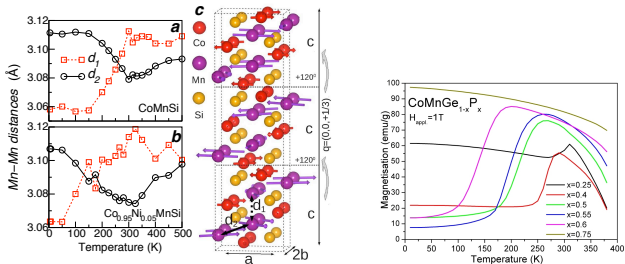
$$a(\vec{m}_0, \vec{B}_t, T_t) = 0, \quad b(\vec{m}_0, \vec{B}_t, T_t) = 0$$

CoMnSi, a metallic helical metamagnet

- Field and T induced transition from low T incommensurate, helical AF state to high magnetisation state.

(A.Barcza, Z.Gercsi, K.S.Knight and K.G.Sandeman, PRL, **104**, 247202, (2010))

- Tricritical point observed.
- Orthorhombic structure. Large magnetoelastic effects.



- Variations, $\text{Co}_{1-x}\text{Ni}_x\text{MnSi}$, $\text{CoMn}_{1-x}\text{Cr}_x\text{Si}$, $\text{CoMnGe}_{1-x}\text{P}_x$

(Z.Gercsi et al., PRB,**83**, 174403, (2011); A.Barcza et al. PRB,**87**, 064410, (2013))

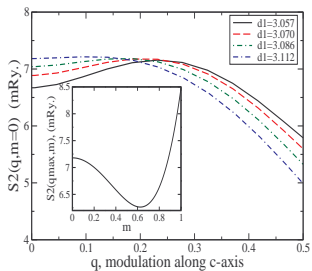
Magnetic states and Disordered Local Moments

- Several solutions $\{\vec{m}_i\}^{(n)}$, find lowest $F(\{\vec{m}_i\}^{(n)}, B \hat{x}, T)$.
- Distorted helix
$$\vec{m}_i^{dh} = (m + \Delta m_{\vec{Q}}) \cos(\vec{Q} \cdot \vec{R}_i) \hat{x} + \Delta m_{\vec{Q}} \sin(\vec{Q} \cdot \vec{R}_i) \hat{y}.$$
- Conical helix: $\vec{m}_i^{ch} = m \hat{x} + \Delta m_{\vec{Q}} (\cos(\vec{Q} \cdot \vec{R}_i) \hat{y} + \sin(\vec{Q} \cdot \vec{R}_i) \hat{z})$.
- Fan state: $\vec{m}_i^{fan} = m \hat{x} + \Delta m_{\vec{Q}} \cos(\vec{Q} \cdot \vec{R}_i) \hat{y}$.
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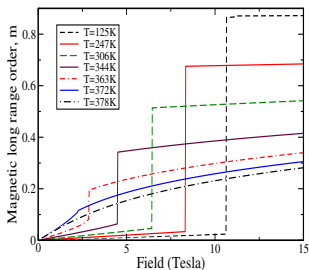
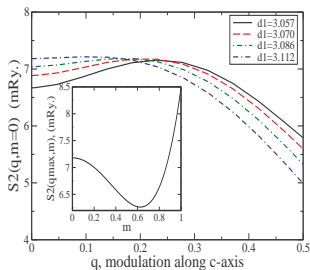
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- FM state: $\vec{m}^{FM} = m \hat{x}$ with $\Delta m_{\vec{Q}} = 0$.
- Relative difference between $F_{distorted-helix}$ and others determines presence or not of first order metamagnetic transition and $\vec{B}_c(T)$.

Metamagnetism in **CoMnSi**

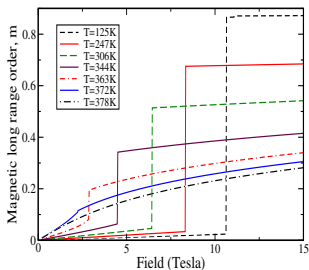
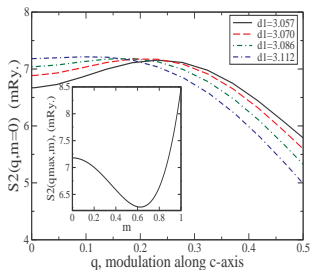


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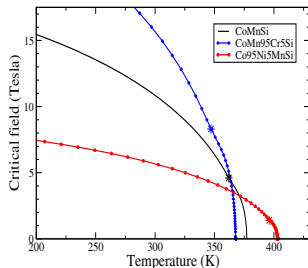
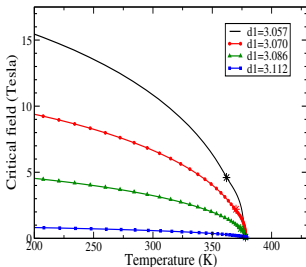
- $S^{(2)}(\vec{Q}, \vec{m} = 0)$, for structures measured in neutron diffraction experiments labelled by the Mn-Mn spacing, d_1 , in Å. The inset shows $S^{(2)}(\vec{Q}_{max}, \vec{m})$ versus FM order parameter m for $d_1 = 3.07$ Å. Spin-orbit coupling and magnetic anisotropic effects very small.
- As d_1 decreases, AFM correlations become stronger than FM ones. $T_N \approx 400\text{K}$. Crucially magnetic interactions weaken as \vec{m} grows. This itinerant electron effective comes from the spin-polarised density induced on the Co sites.

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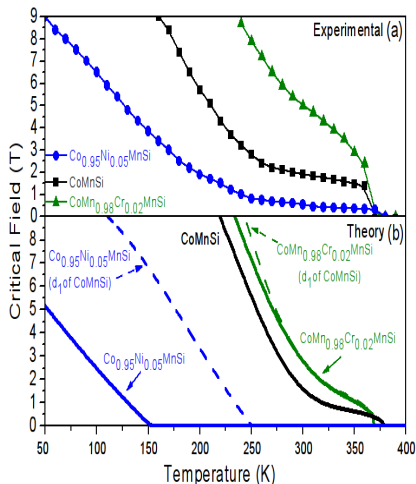
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- m versus applied field for several temperatures for $d_1 = 3.07$ Å. A tricritical point is indicated at 372K, 2 Tesla.

Critical field to switch HAFM to FM/fan state



- The critical field B_c for transition between a helical AFM and a FM state versus T .
- The critical field B_c for $\text{Co}_{0.95}\text{Ni}_{0.05}\text{MnSi}$ (adding electrons) and $\text{CoMn}_{0.95}\text{Cr}_{0.05}\text{Si}$ (removing electrons) with $d_1 = 3.06$ Å.

Test against experiment



Further details: J. B. Staunton et al., PRB 87, 125115(R), (2013).

- Summary:

- Slowly varying **local moments** in sea of 'fast' electrons.
- Theory of metallic magnetism at finite temperatures.
 - MCE of **Gd** around second order FM transition.
 - Application to B2-ordered **FeRh**. First order FM to AF transition. *Strong variation with compositional disorder.*
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- Outlook:

- Find new adaptive magnetic materials.
- Materials modelling tool for magnetic refrigeration materials.
- Nanostructuring magnetic properties.
- Electronic effects, temperature and spintronics.