

Thermodynamic models of the magnetocaloric effect at first order phase transitions

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EU - FP7 Collaboration

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DRREAM
Drastically Reduced Use of Rare Earths in Applications of Magnetocalorics
K. Sandeman (coord.)

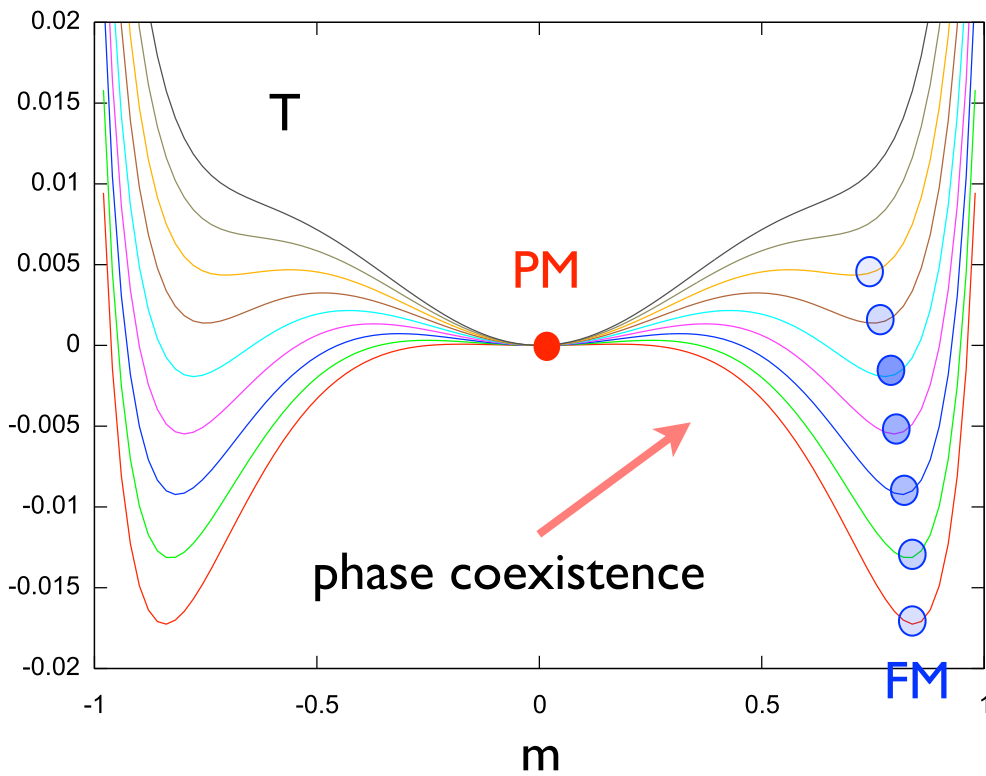
Outline

- 1. Thermodynamic models of MCE**
- 2. Application to $\text{La}(\text{Fe-Mn-Si})_{13}\text{-H}$**
- 3. Conclusions**

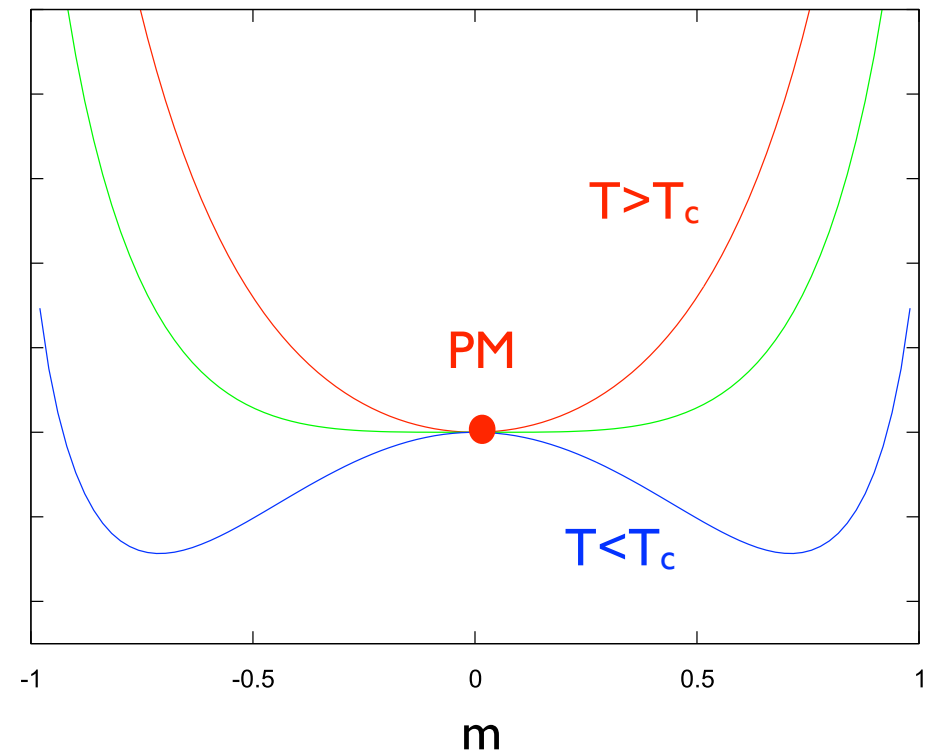
Thermodynamics

free energy $f(M, T)$

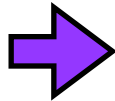
First order transition



Second order transition



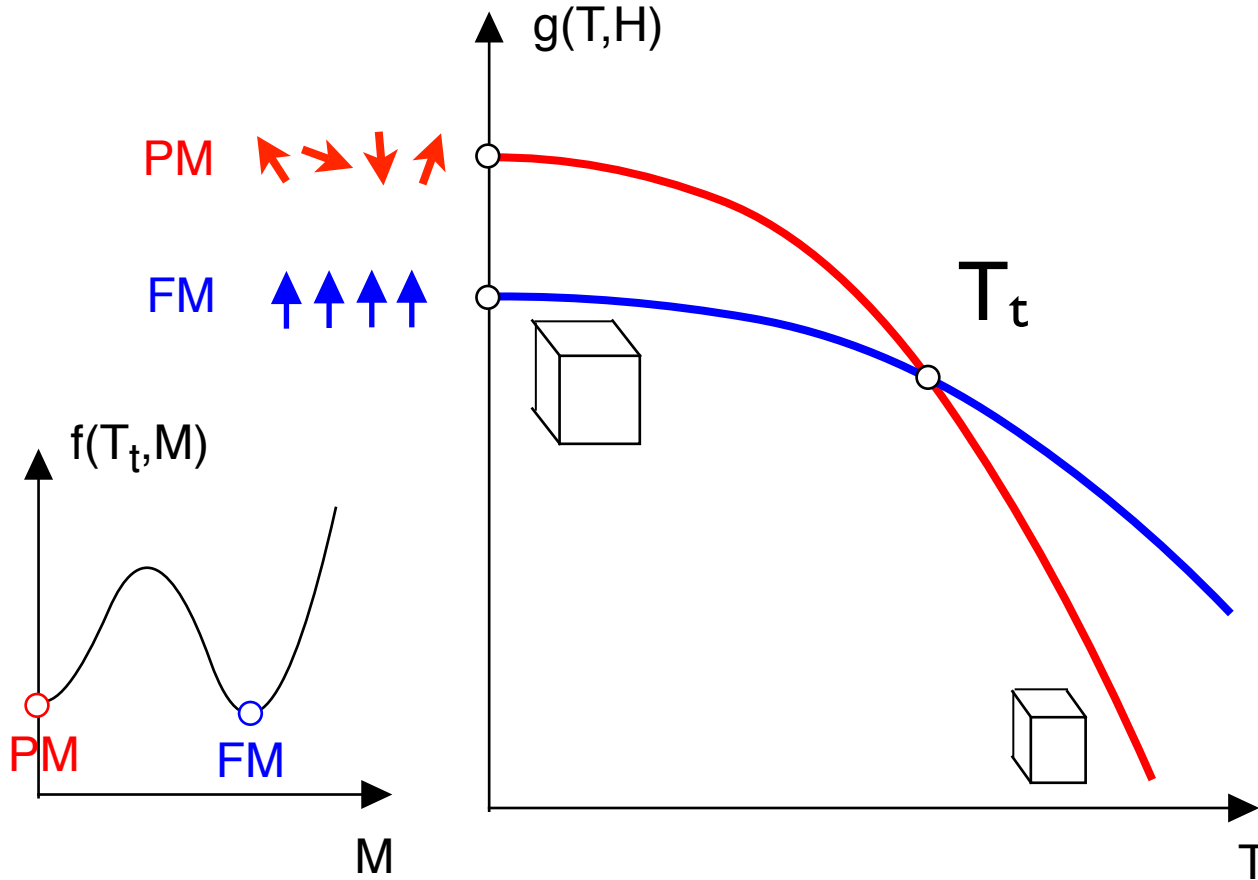
in equilibrium:
transition at $g_{PM} = g_{FM}$



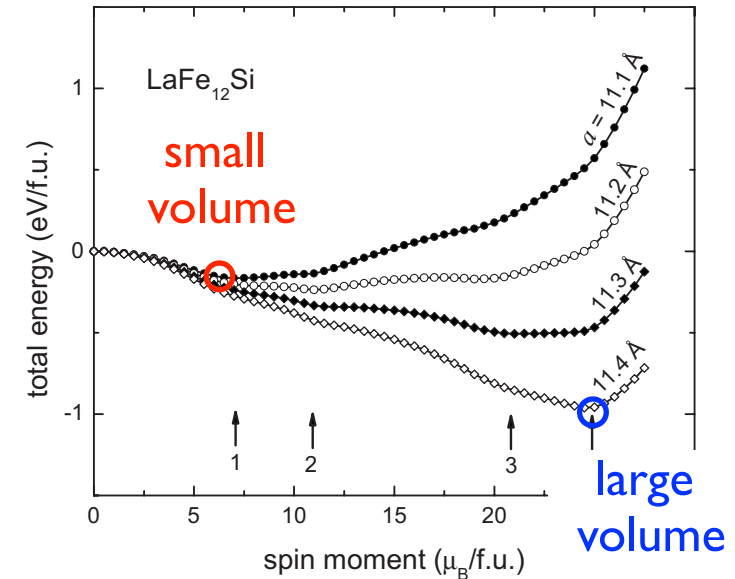
out-of-equilibrium:
hysteresis

Theory

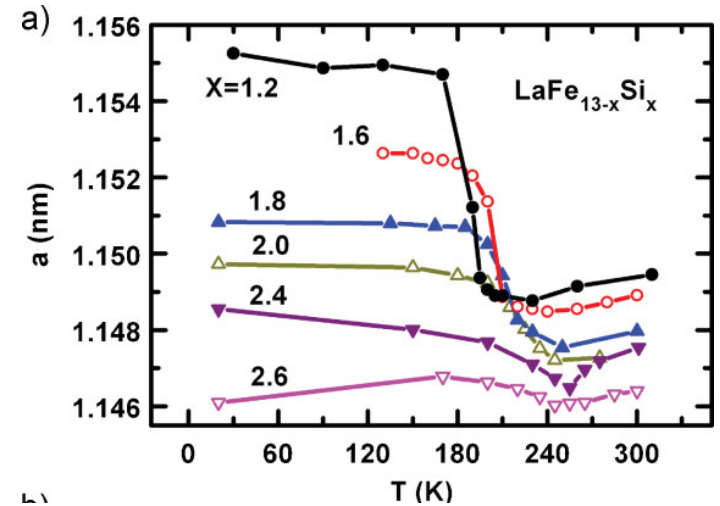
magnetism and lattice



DFT



Kuzmin et al, PRB (2007)



Shen et al, Adv. Mater. (2009)

Free energy

total free energy

$$f_L = -\frac{1}{2}W(\omega)\mu_0 M^2 - T \underline{s_M}(M) + \underline{f_S}(\omega, T)$$

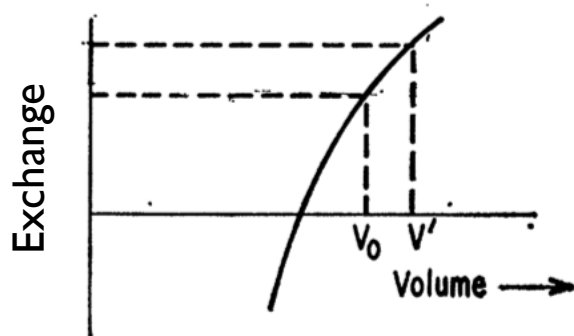
exchange

electronic spin
entropy

all degrees of freedom other than
magnetic

$$K_B N_A \ln(2J+1)$$

lattice free energy



$$\omega = \frac{v - v_0}{v_0}$$

V. Basso, J.Phys. Condens. Matter **23** (2011) 226004.

lattice free energy

$$f_S(\omega, T) = \underbrace{f_{ela}(\omega)}_{\text{elastic}} + \underbrace{f_D(\omega, T)}_{\text{phonons}} + \underbrace{f_{ele}(T)}_{\text{electrons}}$$

- $f_{ela}(\omega) = \frac{v_0}{2\kappa_0} \omega^2 + \mathcal{O}(\omega^3)$

isotropic solid

- $f_D = f(0) + 3nk_B T \left[\ln [1 - \exp(-y)] - \frac{1}{3} \mathcal{D}(y) \right]$

Debye

Debye
temperature

$$y = \frac{T_D}{T}$$

$$T_D(\omega) = T_{D_0}(1 - \gamma\omega)$$

Anharmonic
effects

$$\mathcal{D}(y) = \frac{3}{y^3} \int_0^y \frac{x^3}{\exp(x) - 1} dx$$

Debye function

- $f_{ele} = -\frac{\pi^2}{6} n (k_B T)^2 \mathbf{n}(\epsilon_F)$

small contribution...

lattice free energy

linear elasticity for the lattice around T_0

$$\omega = -\kappa_T p + \alpha_p (T - T_0)$$

$$s_S - s_0 = -v_0 \alpha_p p + b_p (T - T_0)$$

3 free parameters

κ_T - isothermal compressibility

α_p - thermal expansion coefficient

c_v - specific heat ($b_v = c_v / T_0$ - entropy coefficient)

lattice free energy around T_0

$$f_S(\omega, T) = \frac{v_0}{\kappa_T} \frac{\omega^2}{2} - \left[\frac{\alpha_p v_0}{\kappa_T} \omega + s_0 \right] (T - T_0) - b_v \frac{1}{2} (T - T_0)^2$$

term representing the Gruneisen parameter of the Debye theory

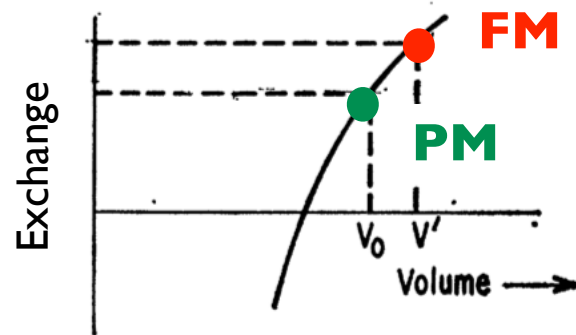
$$b_v = b_p - \alpha_p^2 v_0 / \kappa_T$$

Bean Rodbell model

magneto-volume effects

$$f_L = -\frac{1}{2}W(\omega)\mu_0 M^2 - T s_M(M) + f_S(\omega, T)$$

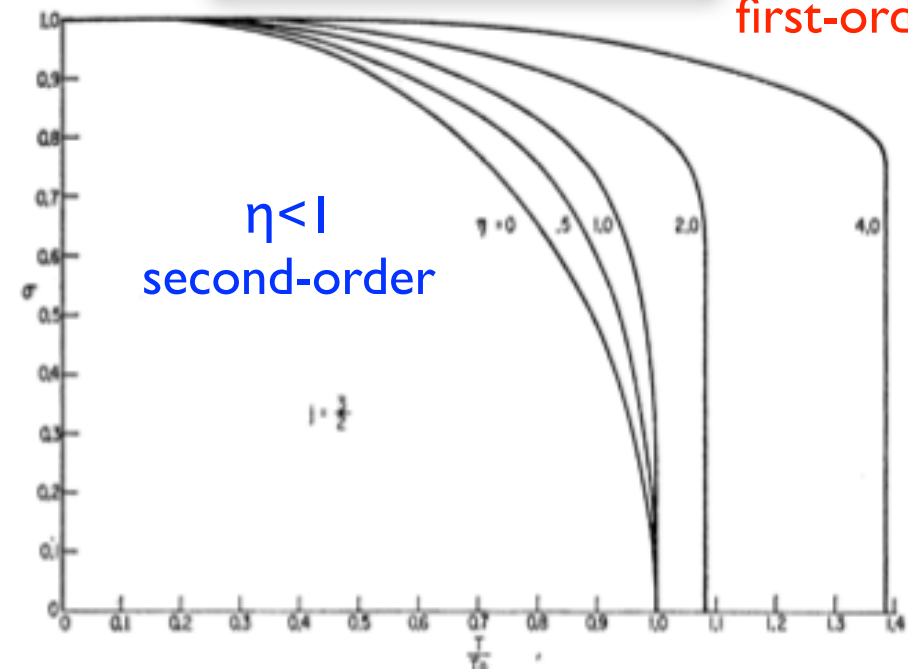
$$\frac{\partial f_L}{\partial \omega} = 0$$



$$W(\omega) = W_0(1 + \beta\omega)$$

$$\eta = \frac{3}{2} \frac{\beta^2 \kappa_T \mu_0 M_0^2 W_0}{v_0}$$

$\eta > 1$
first-order

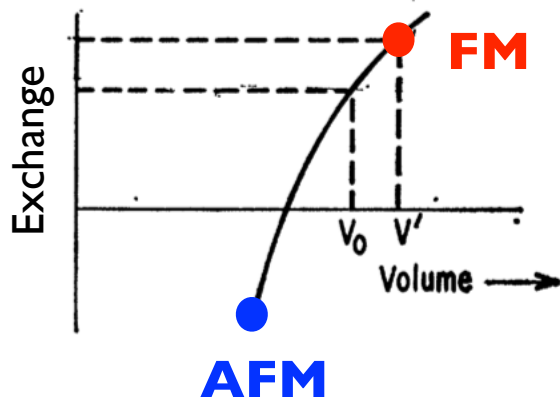


Bean and Rodbell, Phys Rev (1962)

Kittel model

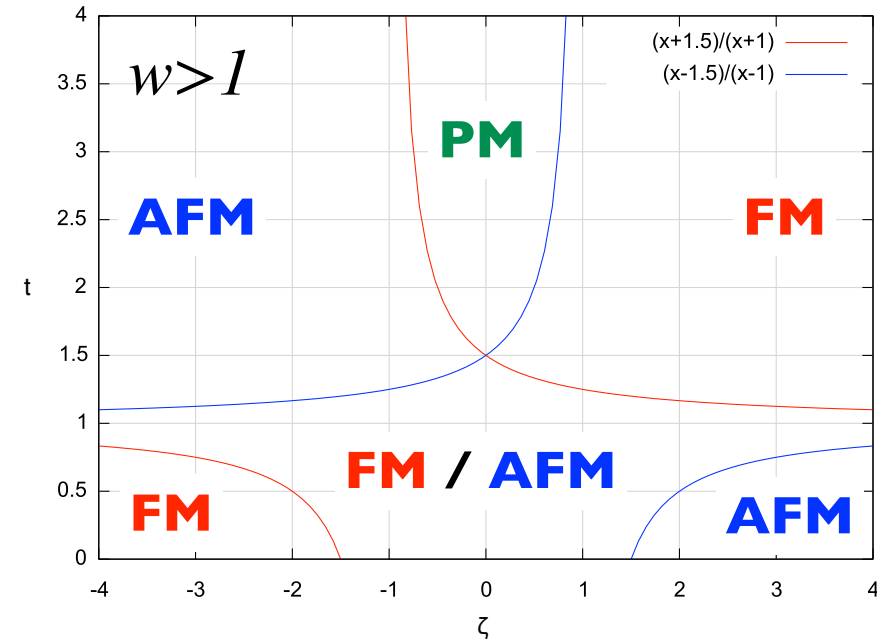
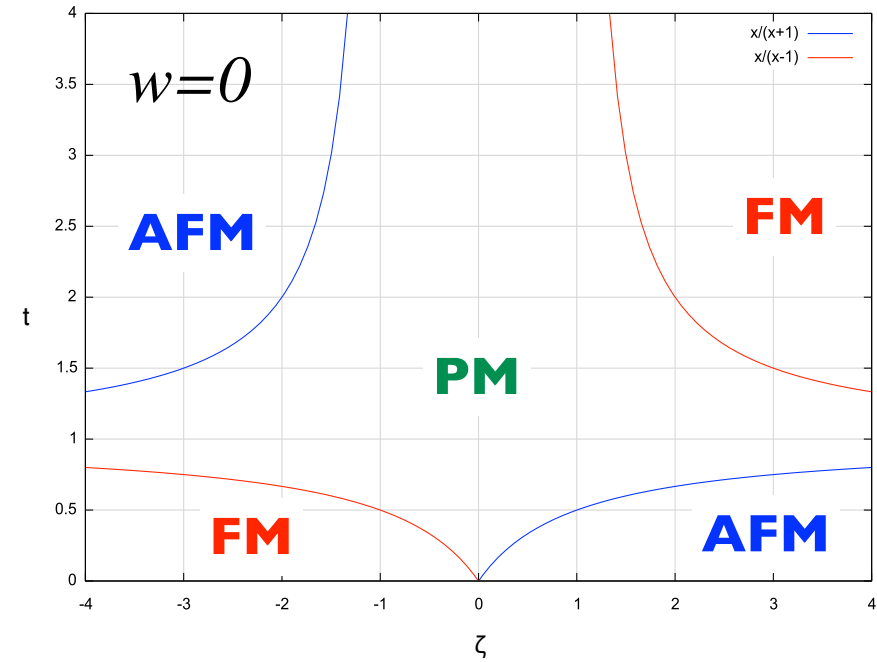
exchange inversion

$$f_{exc} = -2W_{AB}\mu_0 M_A M_B - W\mu_0 M_A^2 - W\mu_0 M_B^2$$



$$W_{AB}(\omega) = W_0(1 + \beta\omega)$$

$$\zeta = \alpha_p \beta T_{c_0}$$



Kittel, Phys. Rev. (1960)

M Piazzzi et al, Physica B (2015)

entropy at the first order transition

magneto-volume effects

total free energy

$$s = - \left. \frac{\partial f_L}{\partial T} \right|_{m, \omega}$$

$$f_L = -\frac{1}{2}W(\omega)\mu_0 M^2 - T s_M(M) + f_S(\omega, T)$$

result

$$s = s_M(m) + s_W(m) + s_S(p, T)$$

magnetic only

$$\bullet \quad s_j(m) = \left[\ln(2j+1) - \frac{1}{a_j} \left(\frac{1}{2}m^2 + \frac{b_j}{4}m^4 + \mathcal{O}(m^6) \right) \right]$$

magneto-elastic

$$\bullet \quad s_W(m) = \frac{nk_B}{2a_J} \zeta m^2 \quad \text{of structural lattice origin}$$

lattice only

$$\bullet \quad s_S - s_{S_0} = -v_0 \alpha_p p + b_p (T - T_0)$$

$$\zeta = \alpha_p \beta T_{c_0}$$

$$\hat{s} = \ln(2J+1) - \frac{1}{2a_J} \left[(1 - \zeta) m^2 + \frac{b_J}{2} m^4 + \mathcal{O}(m^6) \right]$$

for FM to PM

GMCE only for $\zeta < 0$!

values for ζ

$$\zeta = \alpha_p \beta T_{c0}$$

for known MCE materials

Gd₅Si₂Ge₂

$$\beta = -80$$

$$\zeta = -0.6$$

$$\zeta < 0$$

SM+SS

MnAs

$$\beta > 0$$

$$\zeta = 0.25$$

$$\zeta > 0$$

SM-SS

Ni₂MnSn

inverse MCE

$$\beta > 0$$

$$\zeta > 1$$

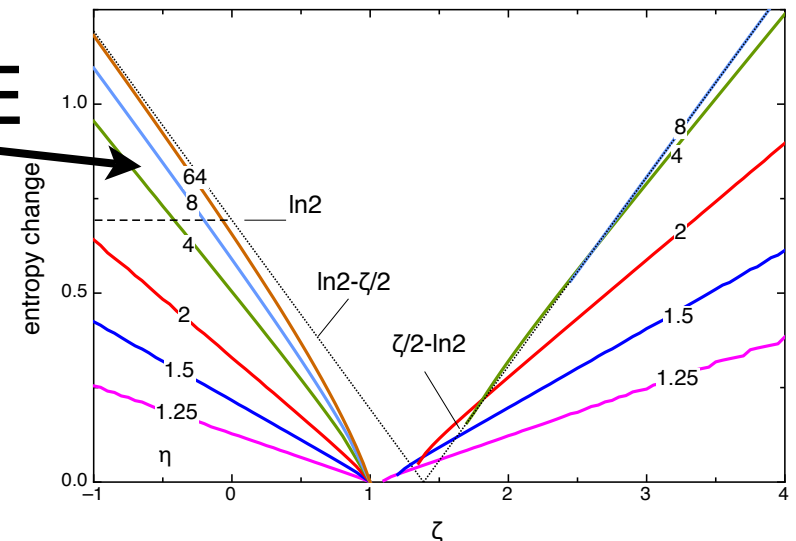
$$\zeta > 1$$

-SM+SS

prediction

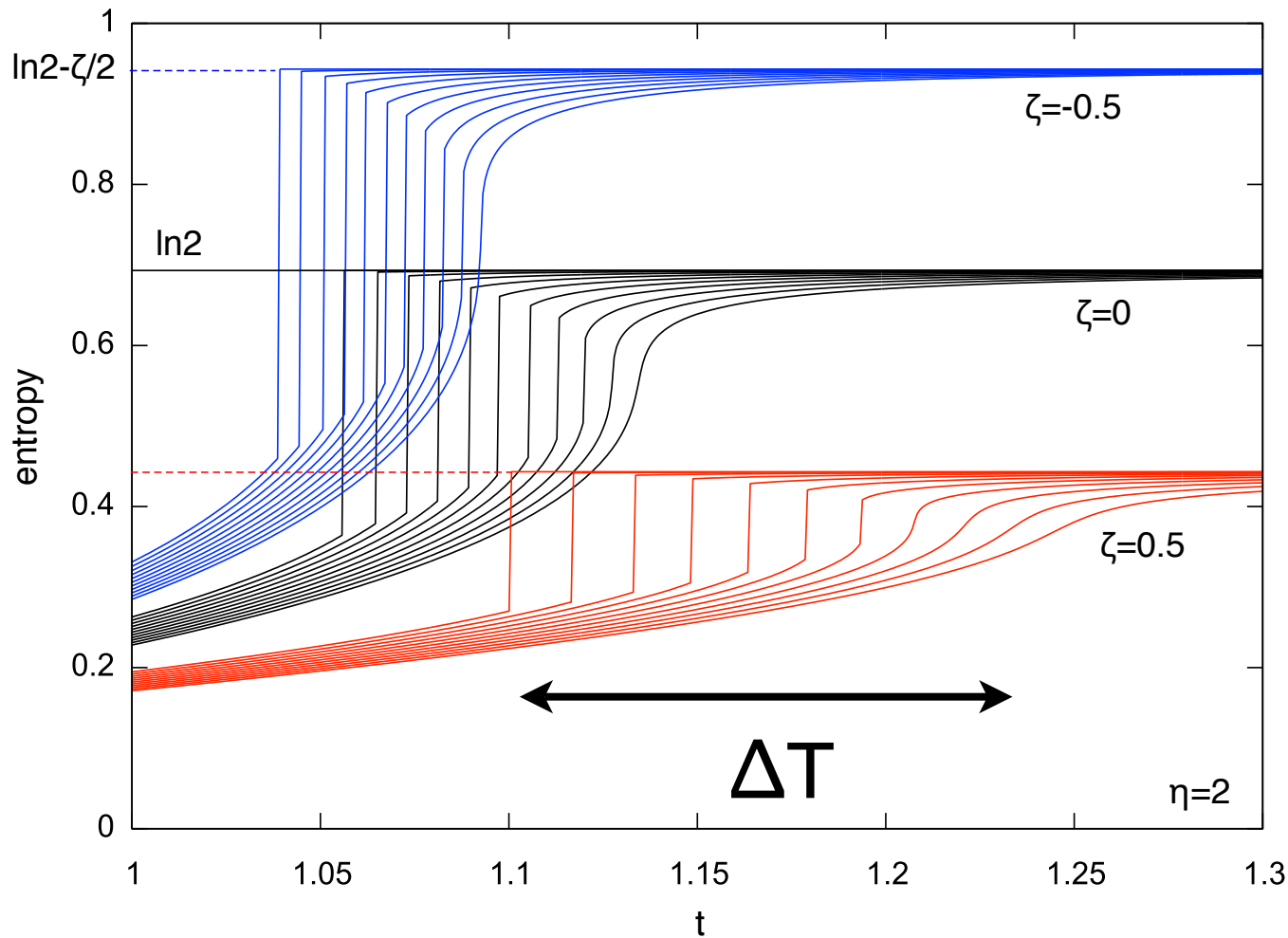
GMCE in Mn alloys ($\beta > 0$) but if thermal expansion is negative

GMCE



entropy at FM-PM first-order transition

magneto-volume effects



GMCE

$\zeta < 0$
SM+SS

$\zeta > 0$
SM-SS

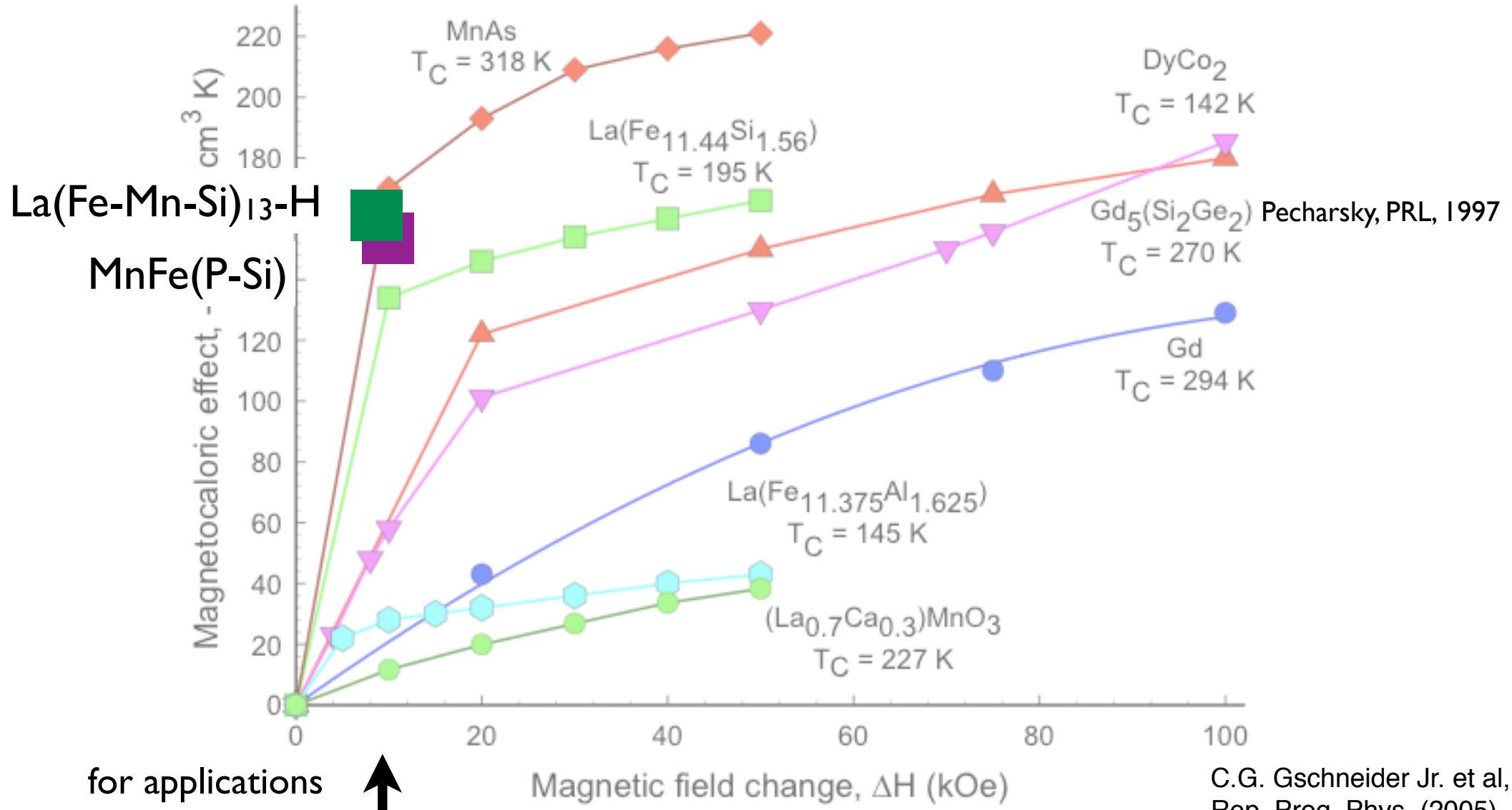
example for spin $j=1/2$

V. Basso, J.Phys. Condens. Matter **23** (2011) 226004.

Magnetic cooling

new materials

with large MCE

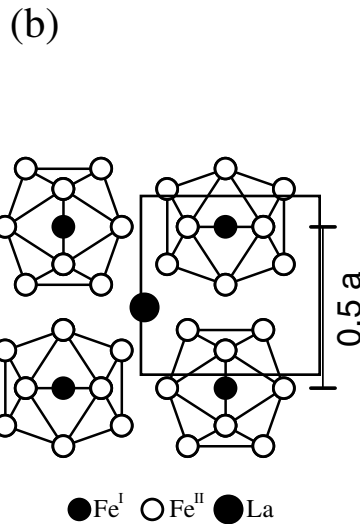
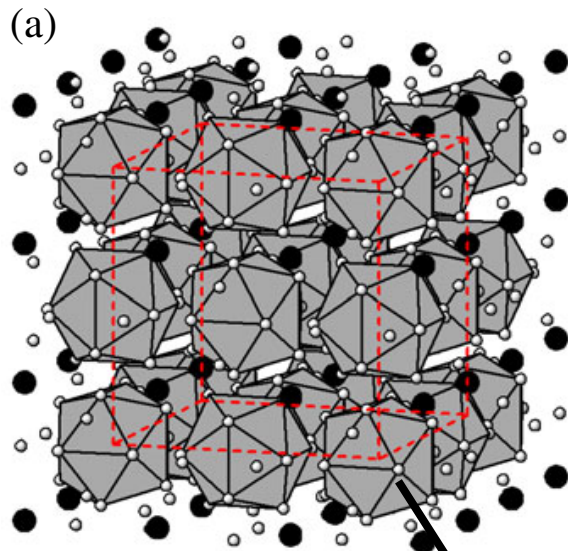


La(Fe-Mn-Si)₁₃H_{1.65}

samples prepared by VAC

combine:

- substitution of Fe with Mn
- hydrogenation at saturation

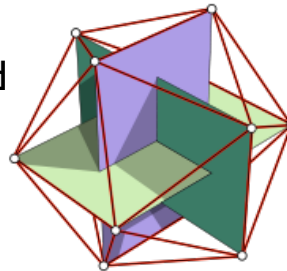


regular icosahedron
(12 vertex)

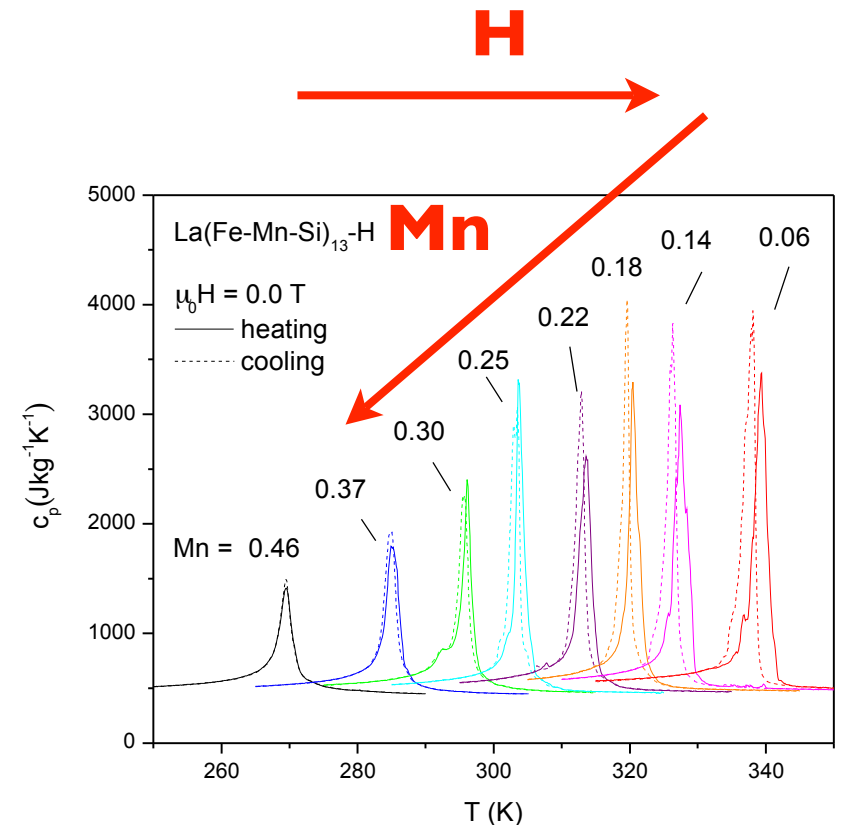
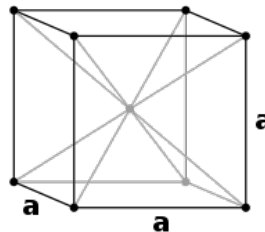
Fe (I) 1.54 μ_B
bonds:
- 12 Fe (I)

Fe (II) 2.16 μ_B
bonds:
- 1 Fe (I) short bond
- 5 inter-cluster Fe(II) short bond
- 2 intra-cluster Fe(II): long bond

fcc-like coord

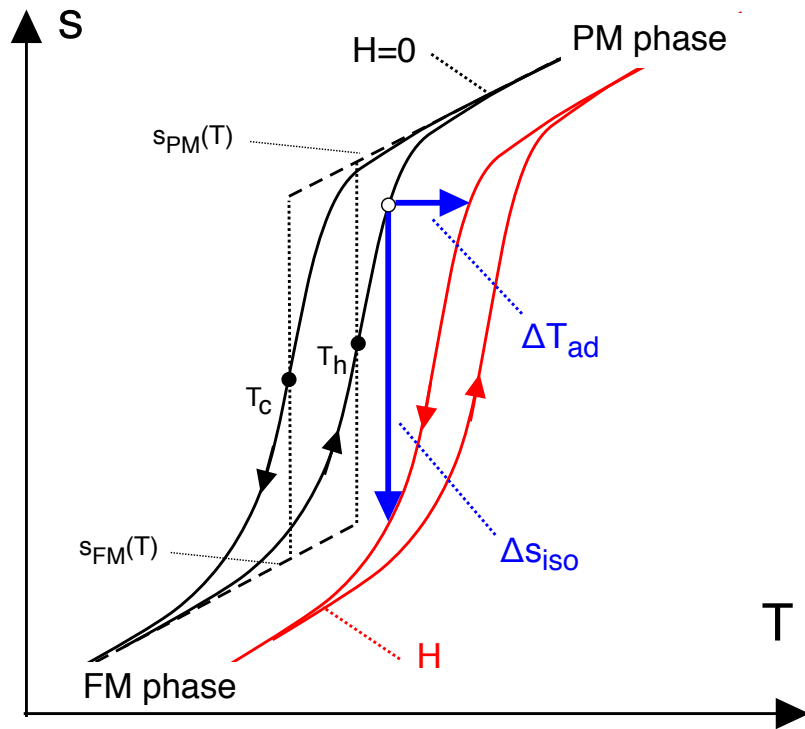


bcc-like coord

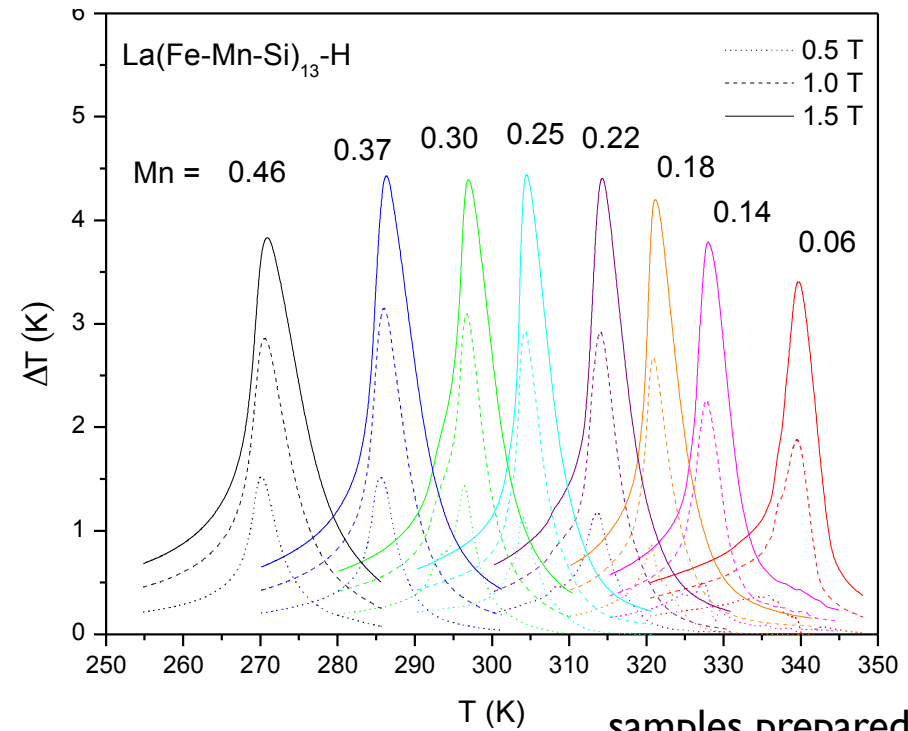
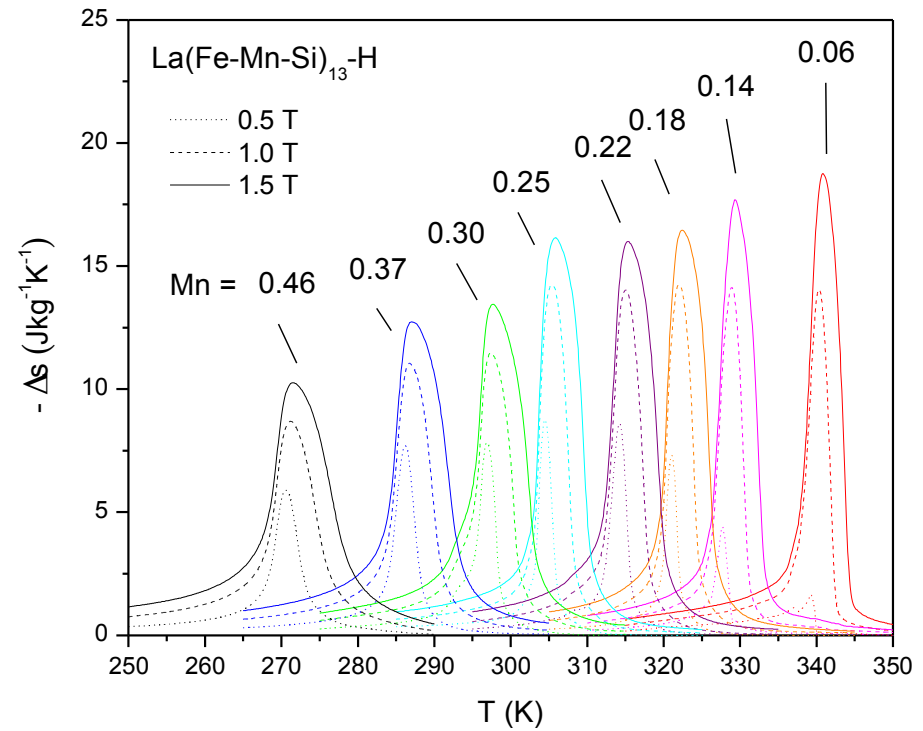


Fujita et al., PRB 2003

La(Fe-Mn-Si)₁₃H_{1.65}

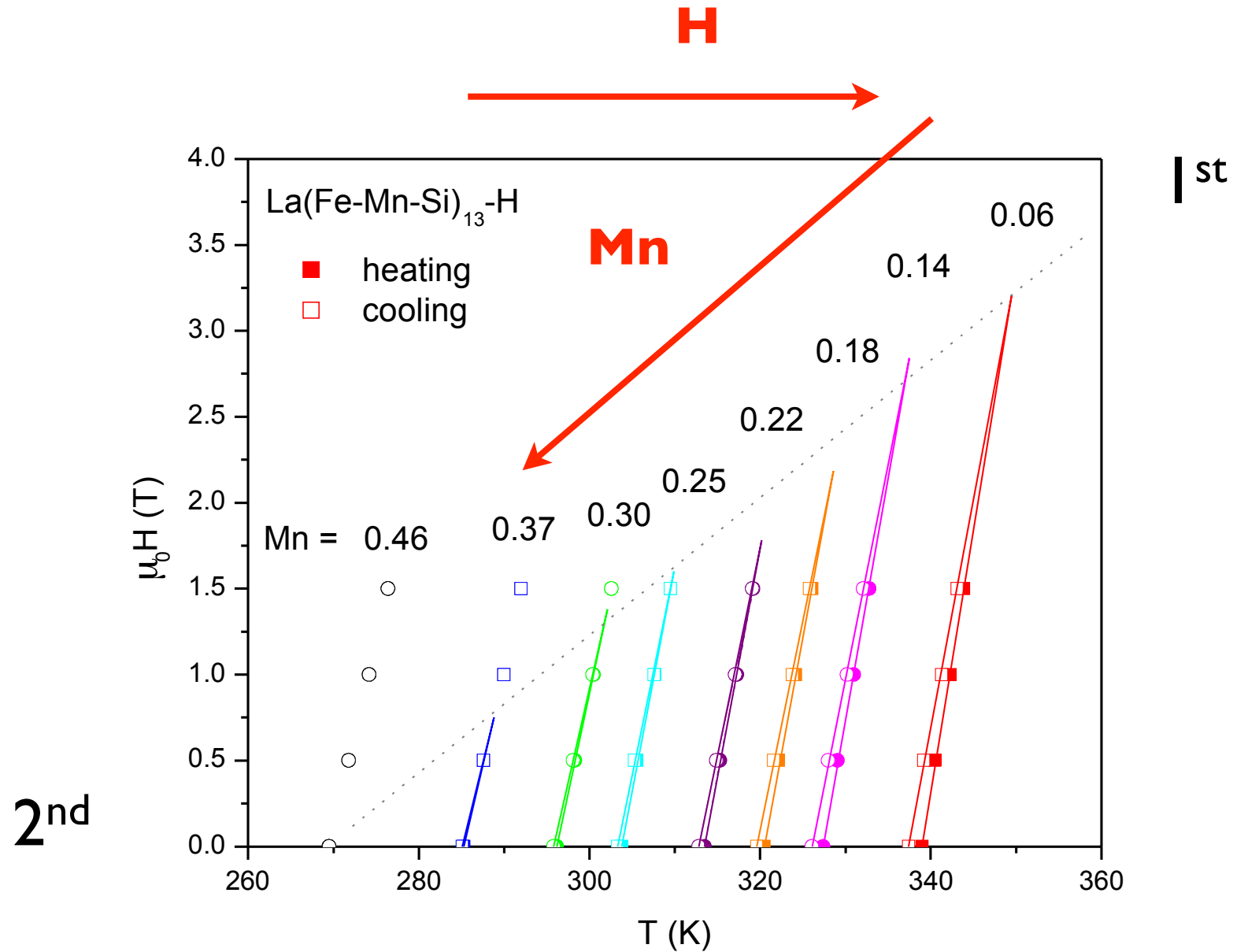


temperature hysteresis limits ΔT_{ad}

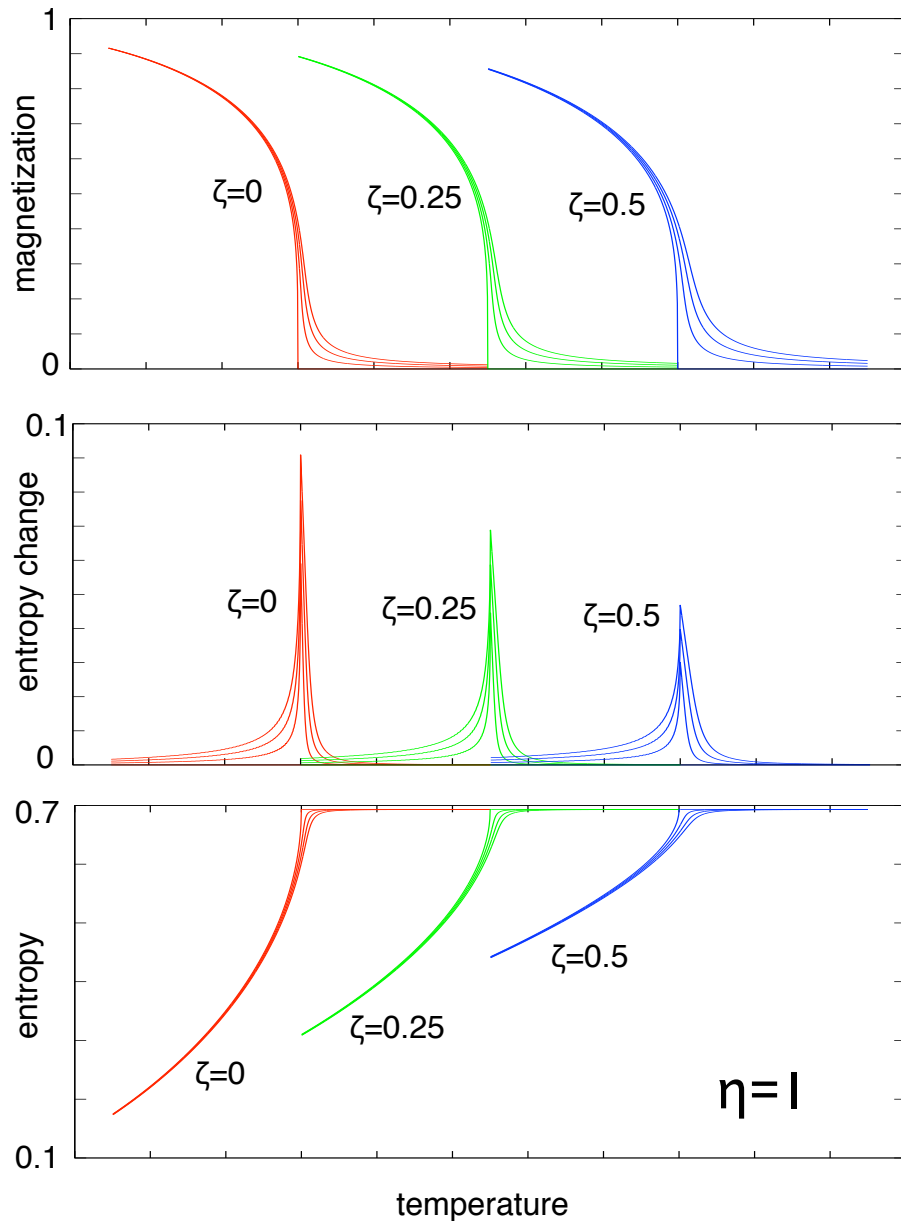


samples prepared by VAC

La(Fe-Mn-Si)₁₃H_{1.65}



La(Fe-Mn-Si)₁₃H_{1.65} model

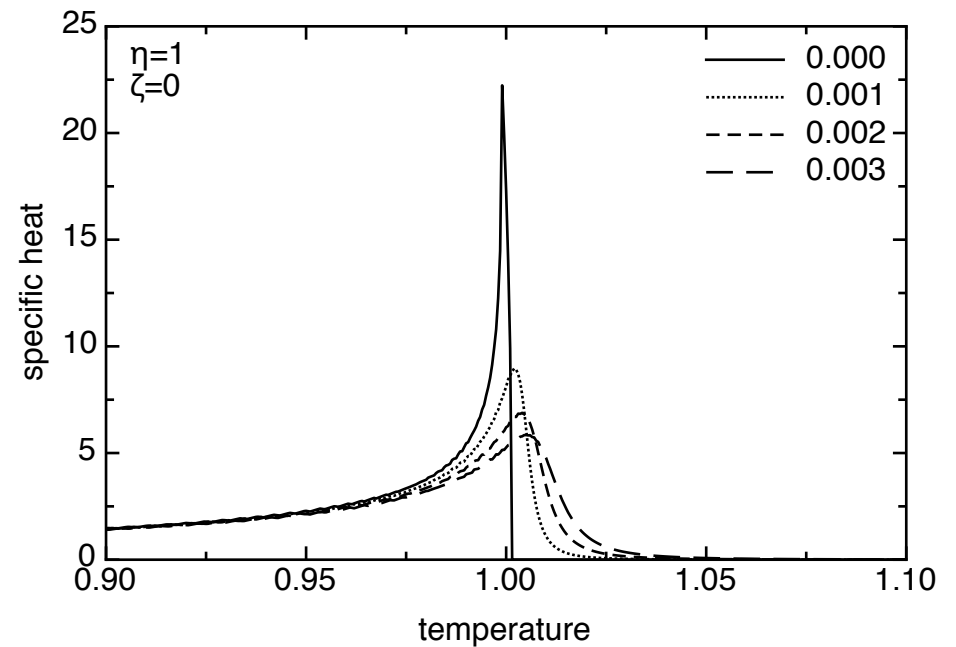


$$\zeta = \alpha_p \beta T_0$$

lattice entropy sums or subtracts

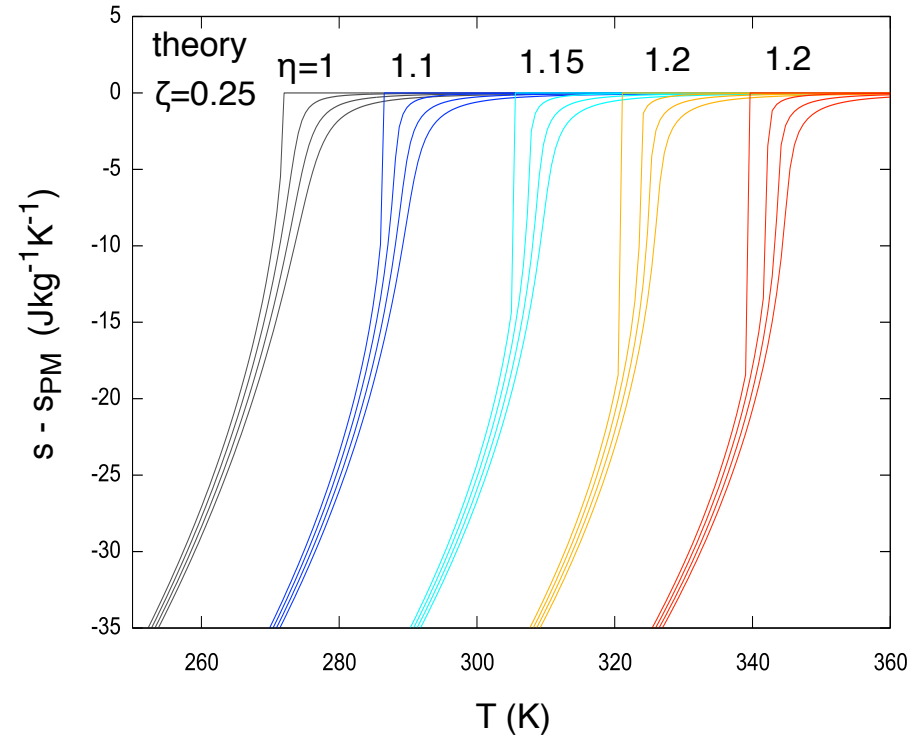
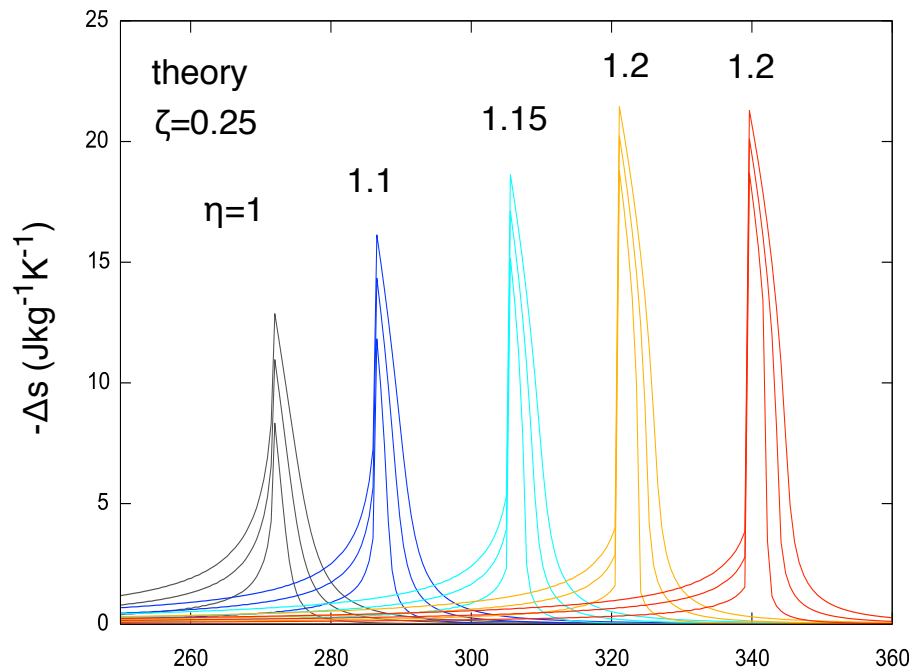
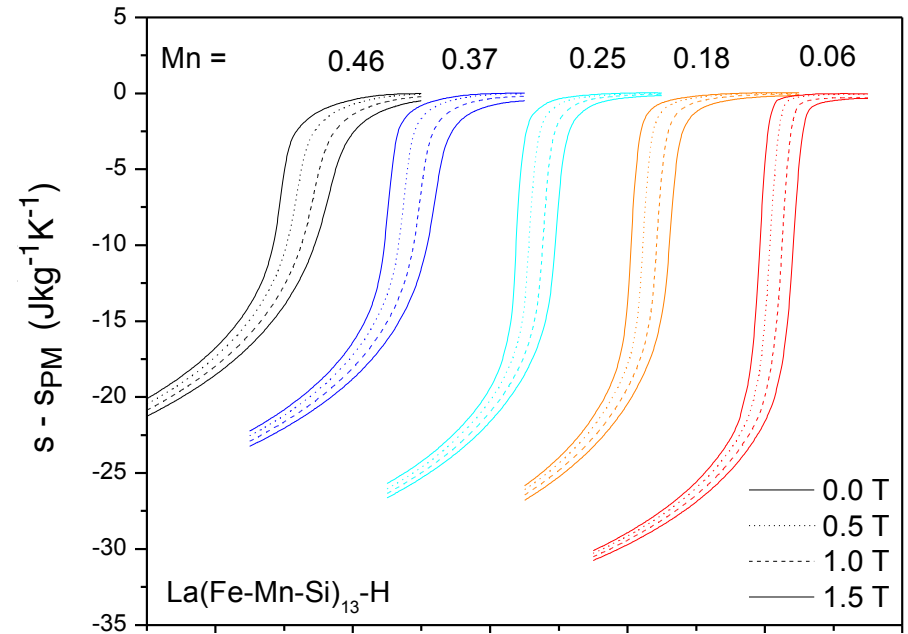
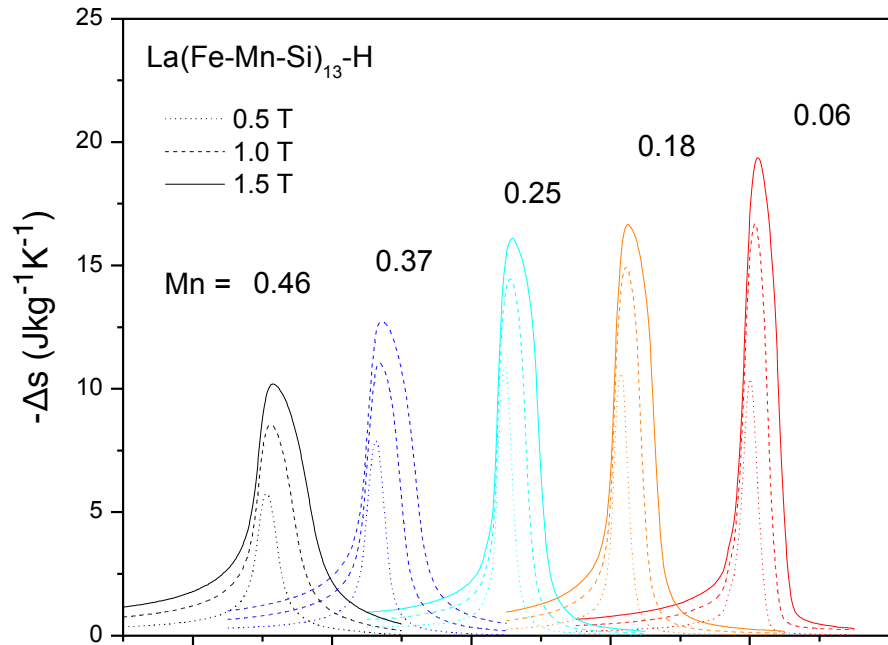
$$\eta = \frac{3}{2} \frac{\beta^2 \kappa_T n k_B}{a_J v_0} T_0$$

first order/second order



M Piazza et al , JMMM (2015)

La(Fe-Mn-Si)₁₃H_{1.65} model



M Piazzzi et al , JMMM (2015)

Conclusions

La(Fe-Mn-Si)₁₃H_{1.65} model

Fe	Mn	Si	T_0 (K)	$\mu_0 H_0$ (T)	η
11.22	0.46	1.32	272	405	1.0
11.33	0.37	1.30	286	426	1.1
11.47	0.25	1.28	305	455	1.15
11.60	0.18	1.22	320.5	478	1.2
11.76	0.06	1.18	339	505	1.2

$$\eta = \frac{3}{2} \frac{\beta^2 \kappa_T n k_B}{a_J v_0} T_0$$

first order/second order

$$\zeta = \alpha_p \beta T_0$$

lattice entropy
sums or subtracts

$$\kappa_T = 8.6 \cdot 10^{-12} \text{ Pa}^{-1}$$

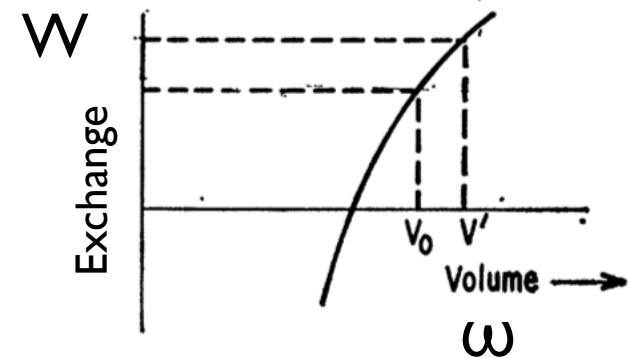
$$\alpha_p \simeq 5 \cdot 10^{-5} \text{ K}^{-1}$$

$$\zeta = 0.25$$

- Mn influences the exchange W_0 (y_{Mn})
- β (magnetoelasticity) depends on the 1:13 structure

$$\beta \simeq 15$$

$$W = W_0(1 + \beta \omega)$$



Perspectives

magnetism and lattice

- Use SDTF results to predict new transitions

$f(\omega, M)$ average magnetization vs local moment

- B. L. Gyorffy, et al. J. Phys. F: Met. Phys., 15:1337 (1985). - *Disordered Local Moment*
- J. B. Staunton et al, Phys. Rev. B, 89 :054427 (2014). - *FeRh*.

- M. Piazzi, et al Physics Procedia (ICM 2015) - *first tests*

.....

Other terms with explicit dependence on M

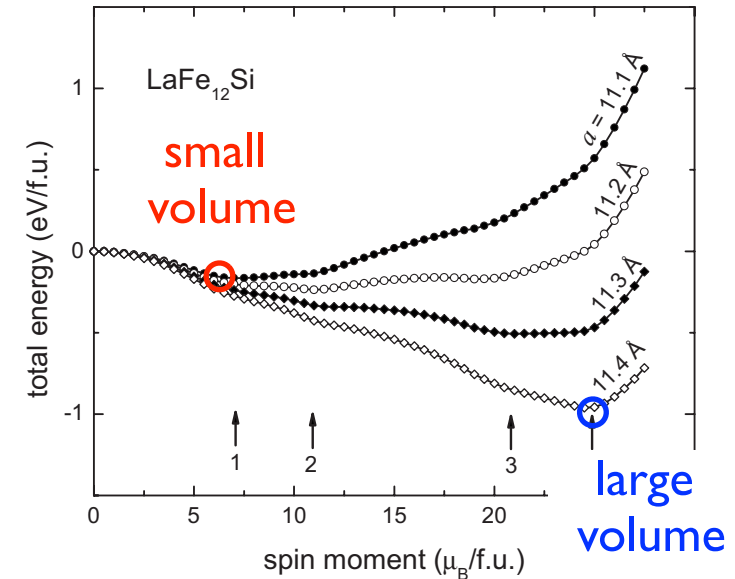
- Structural contribution (phonons)

Gruner et al. PRL 2015 where is spin entropy?

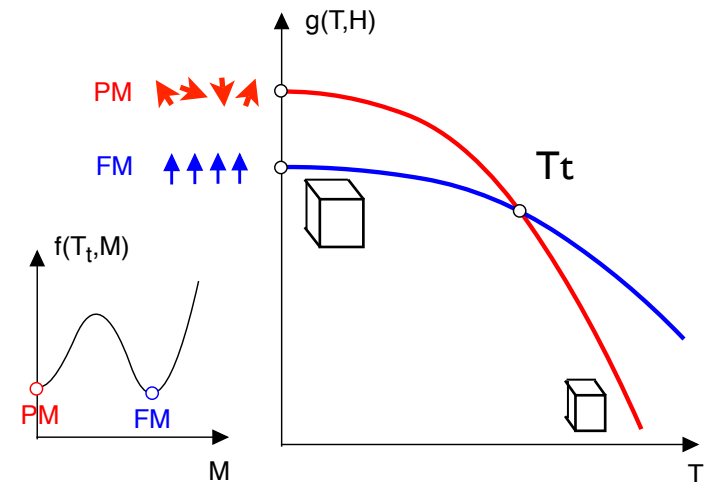
- Electronic contributions (Fermi level)

A. Fujita et al PRB (2002) - *IEM theory*

DFT



Kuzmin et al, PRB (2007)



Acknowledgements

- **EC** 7th Framework Programme
Project **SSEEC** Solid State Energy Efficient Cooling 2008/2011



- **BASF** future business



- **EC** 7th Framework Programme
Project **DRREAM** Drastically reduced rare earth use in applications of magnetocalorics. 2013/2015



7th International Conference on Magnetic Refrigeration at Room Temperature



Thermag VII



Torino, Italy, September 11-14, 2016

www.thermag2016.com

Conference topics:

- **Materials:** alloys and compound, physics, shaping and preparation
- **Cooling systems:** design and numerical aspects, experimental tests.
- **Novel aspects and future perspectives:** ferromagnetic materials, thermal switches, etc



Thanks for your attention