



Pressure effects on the thermal hysteresis in $\text{MnFe}(\text{P},\text{Ge})$ compounds

O. Tegus

**Department Radiation, Radionuclides & Reactors
Faculty of Applied Sciences
Delft University of Technology**

**Inner Mongolia Key Lab
for Physics and Chemistry of Functional Materials
Hohhot, China**

Delft days on magnetocalorics, 31 Oct. 2008

■ Introduction

Experimental results

■ Theoretical model and Experimental results

■ Conclusions

Pressure



Tuning of molar volume



Modification of intermolecular interactions



Destabilization of intermolecular bonds



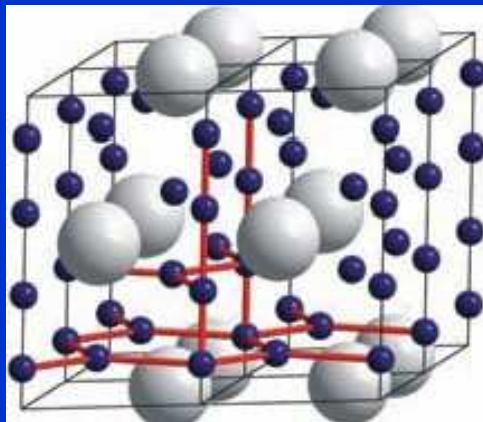
Phase transitions

Examples: Pressure effect

Iron becomes superconductor(<2 K): 10-30 GPa

Hydrogen becomes metal: 400 GPa
(static pressure 1 Matm, dynamic pressure 10 Matm)

Silane becomes Superconductor: 50 GPa



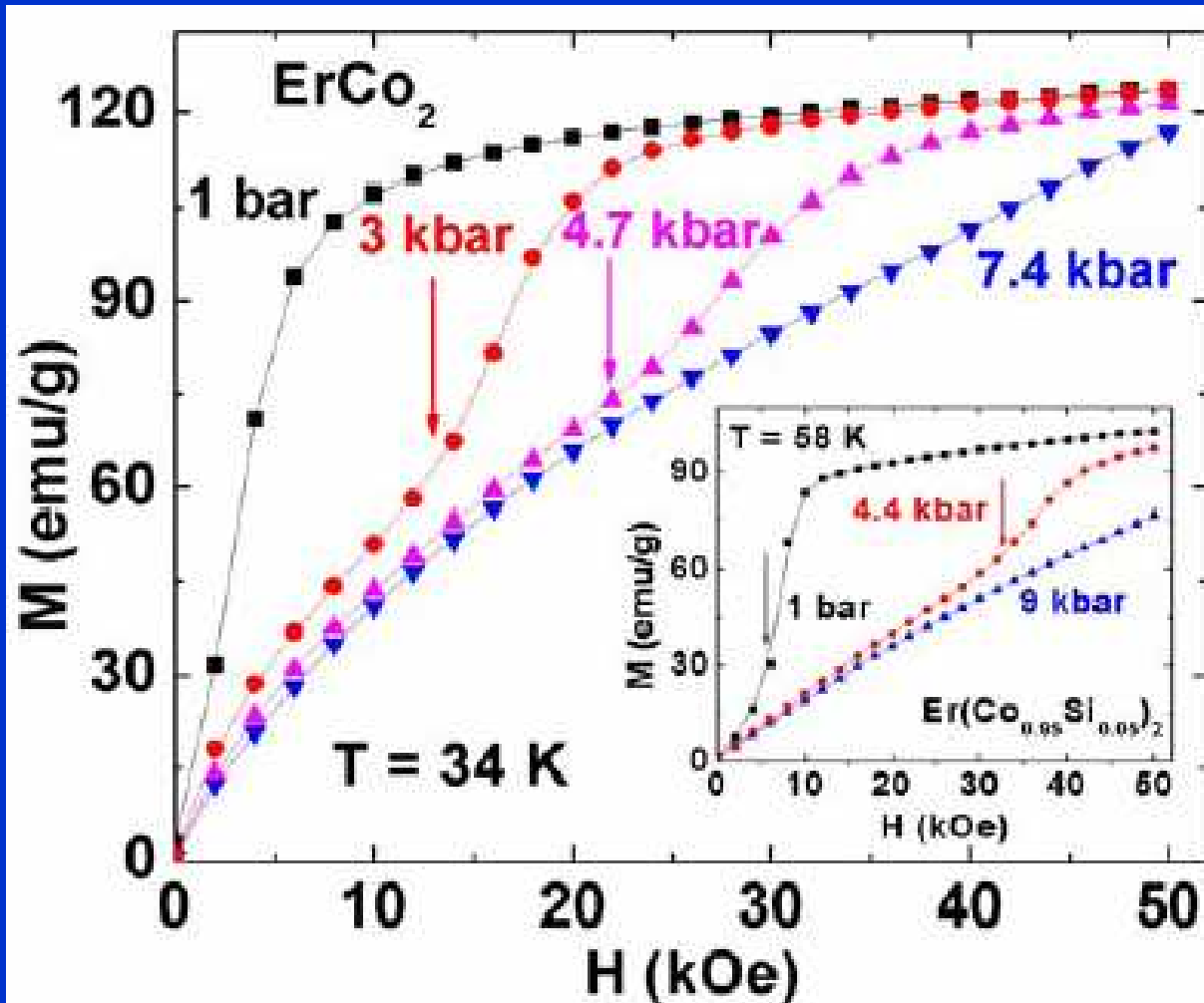
Effects of pressure on magnetism

Pressure effects on magnetization

ErCo_2

1st order phase transition

$T_c = 32 \text{ K}$

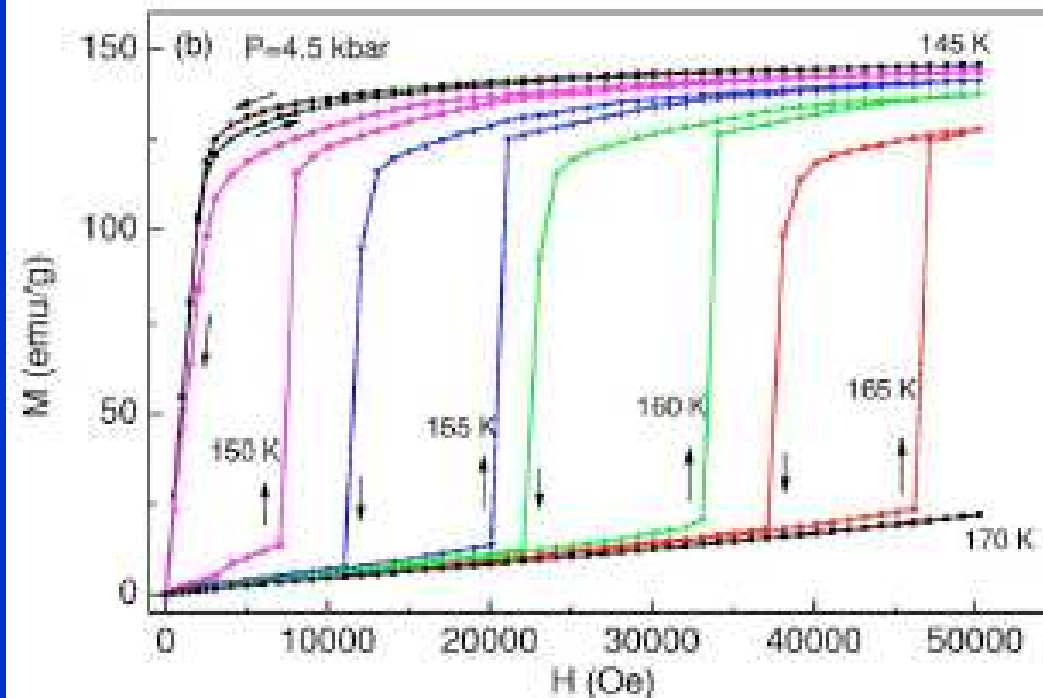
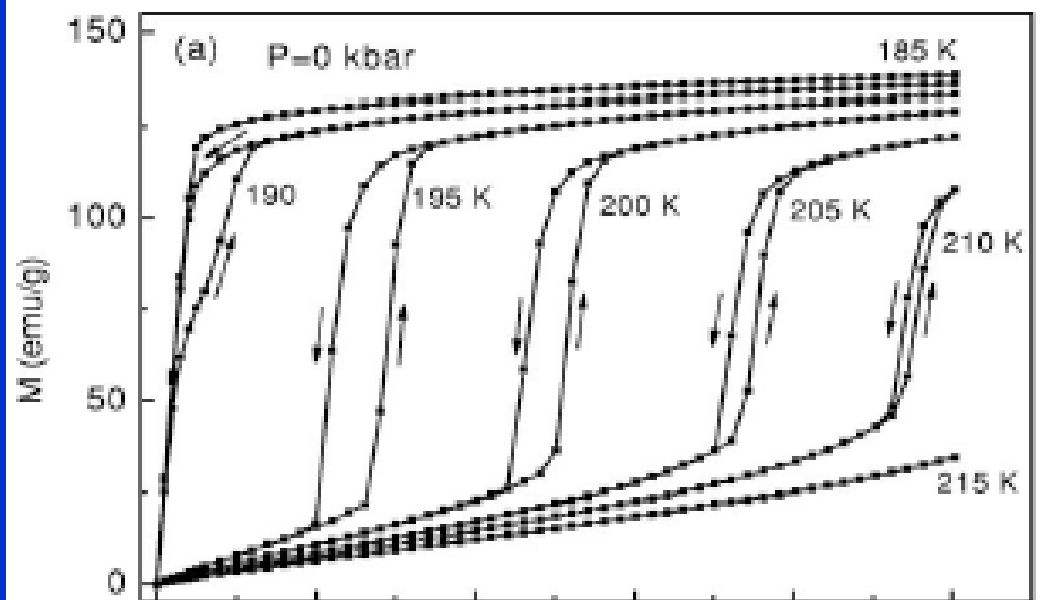




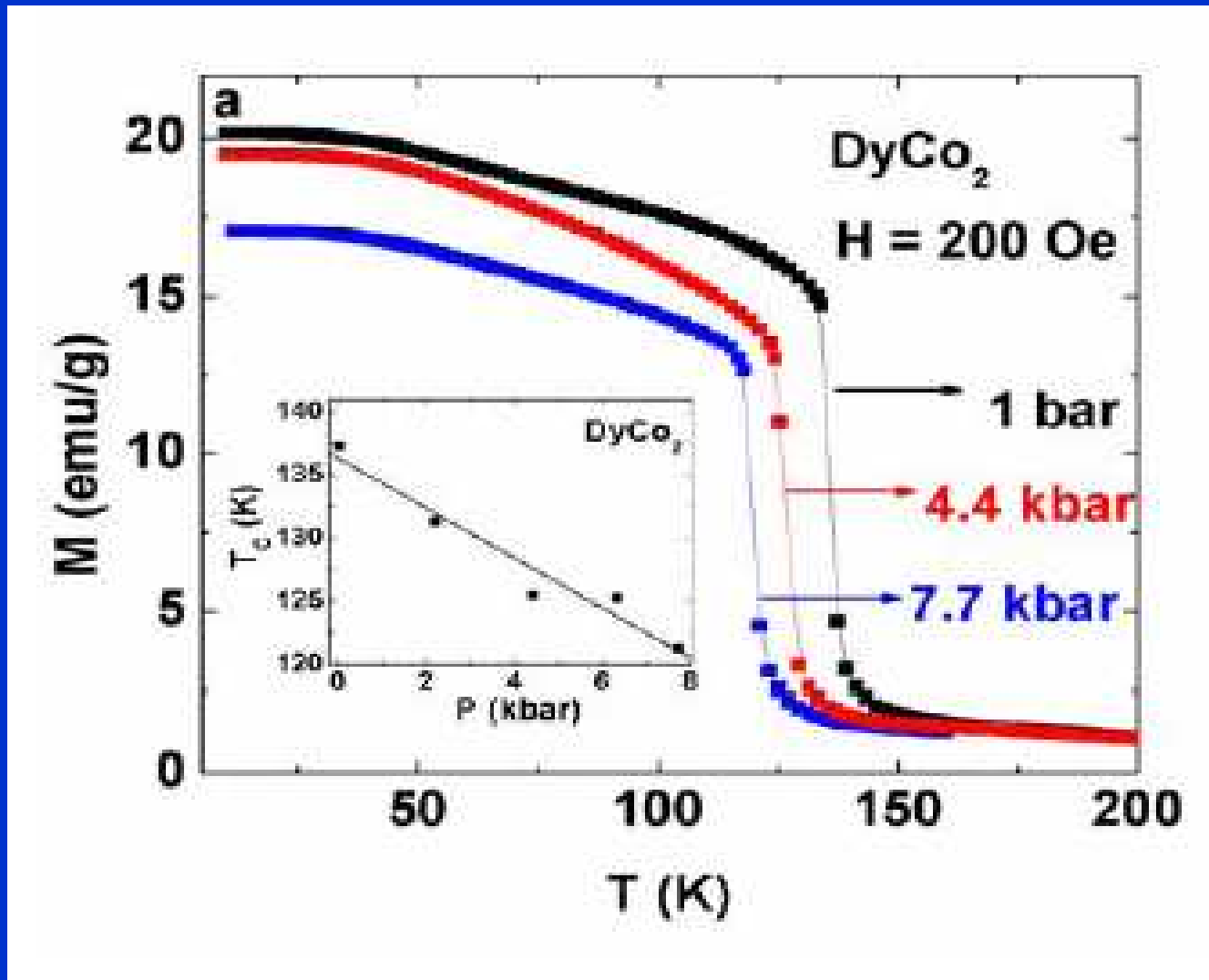
$P = 0$ kbar

$P = 4.5$ kbar

1st order phase transition



Pressure effects on exchange interactions

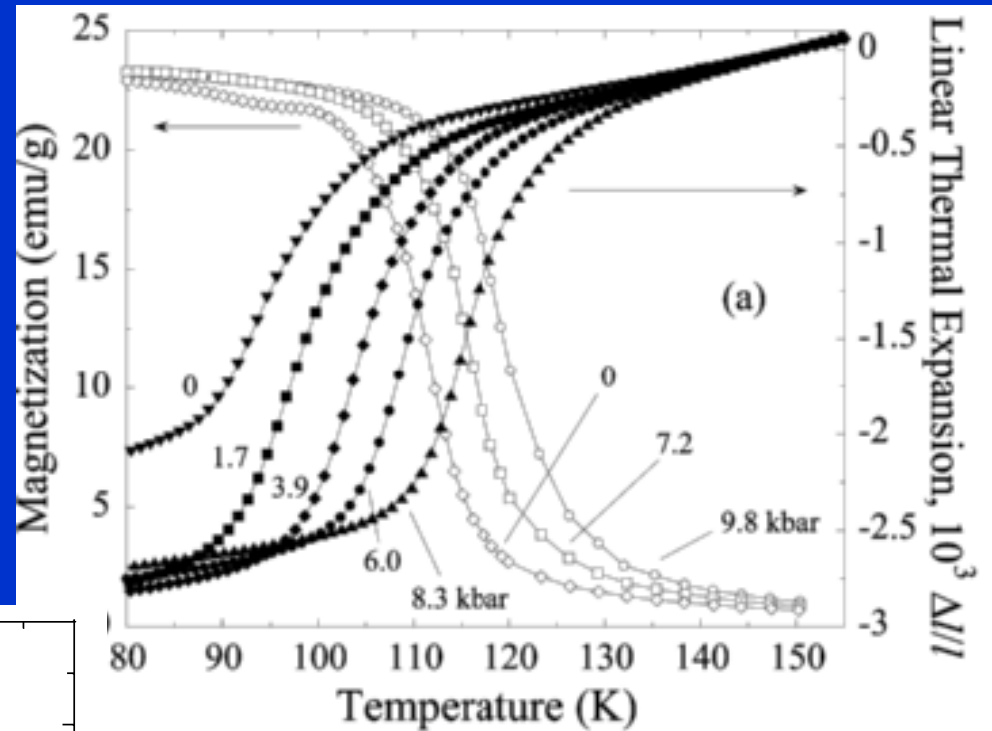
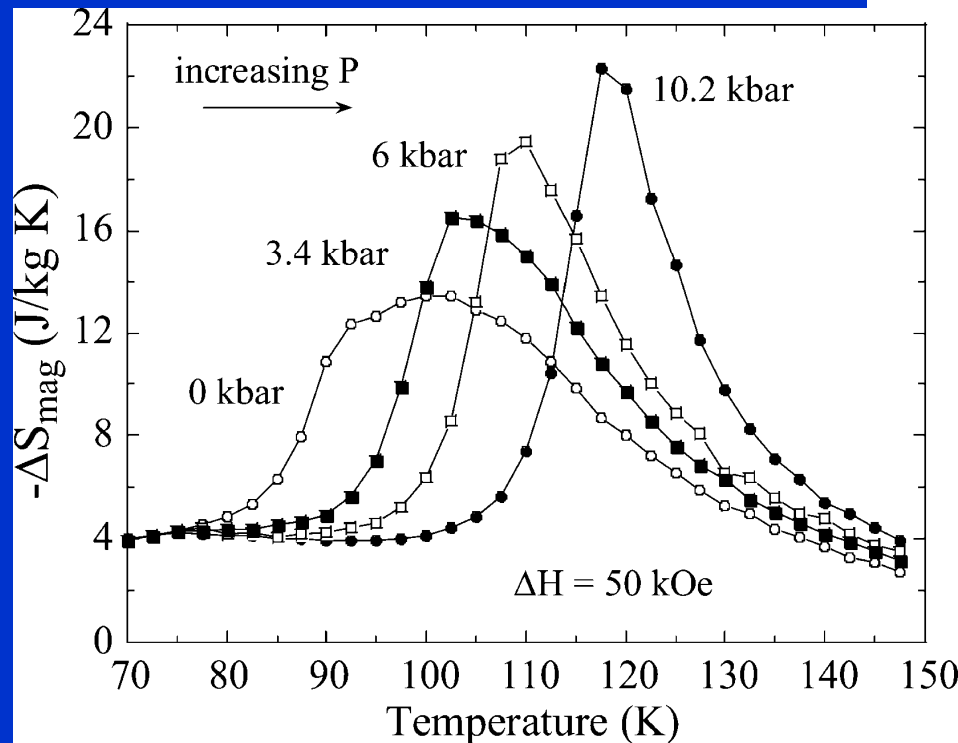


T_h decreases with increasing p

Tb₅Ge₂Si₂

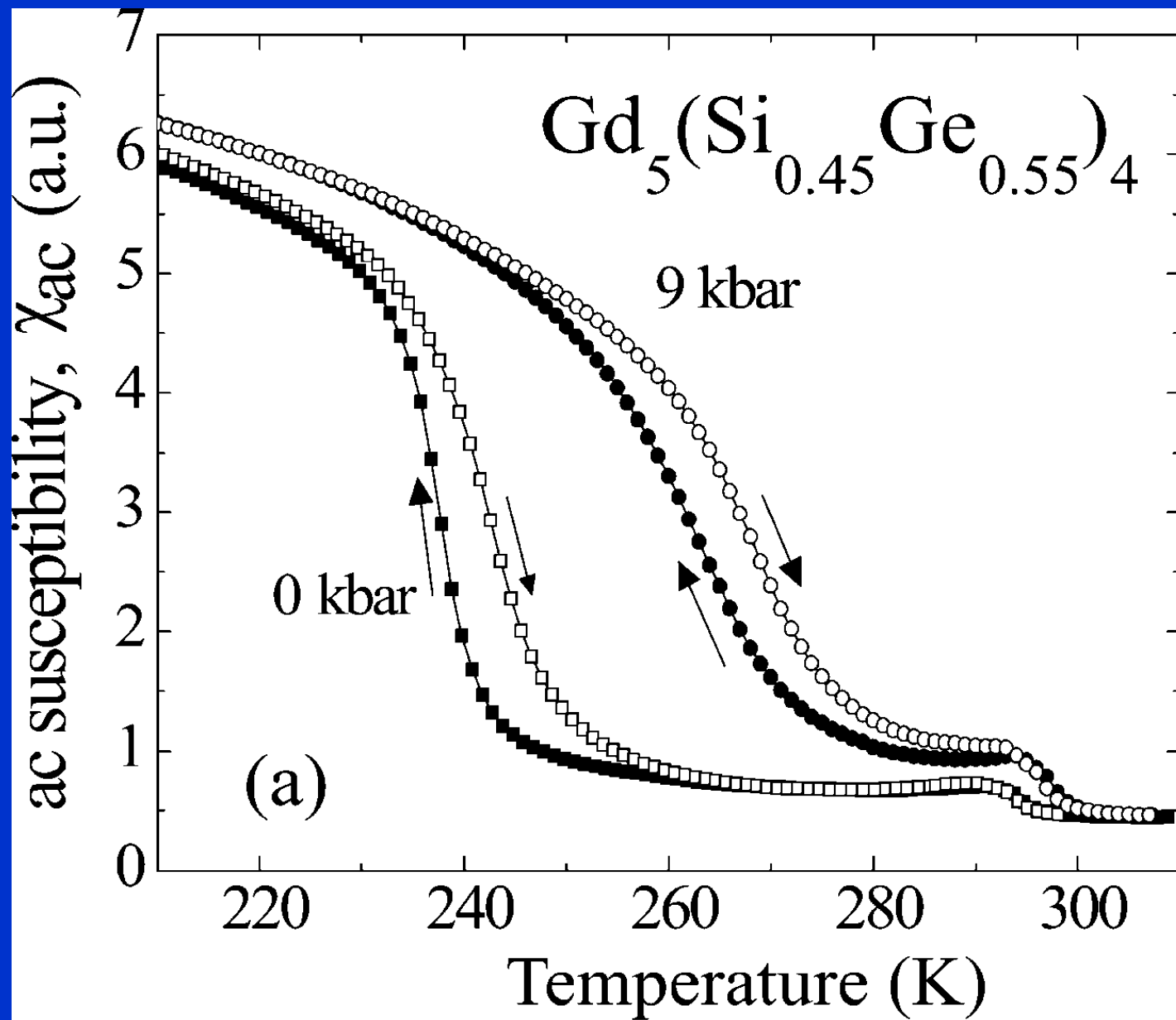
**1st order phase
Transition**

T_c = 105 K

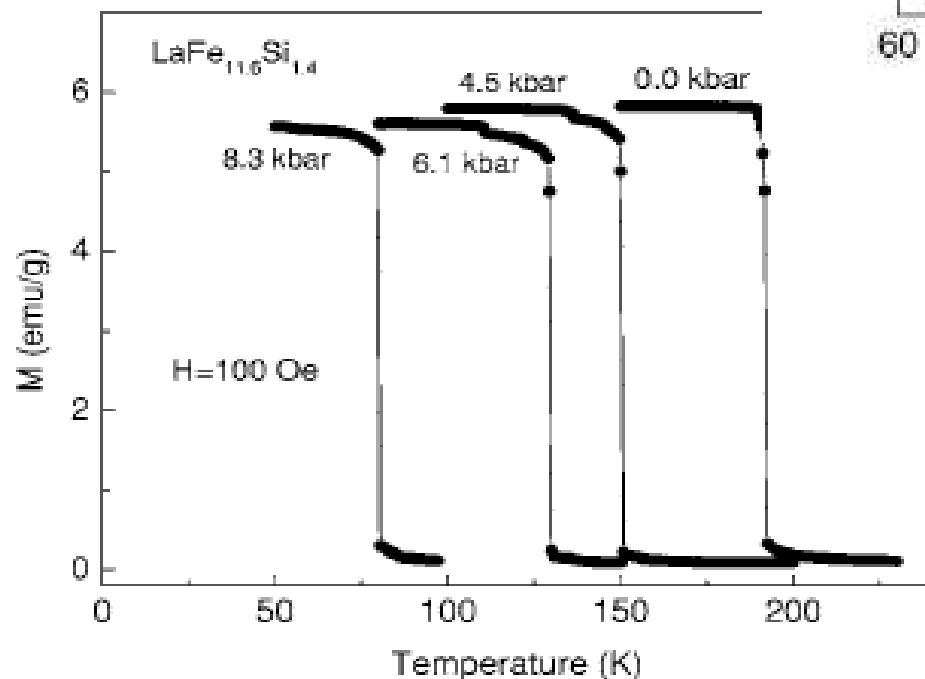
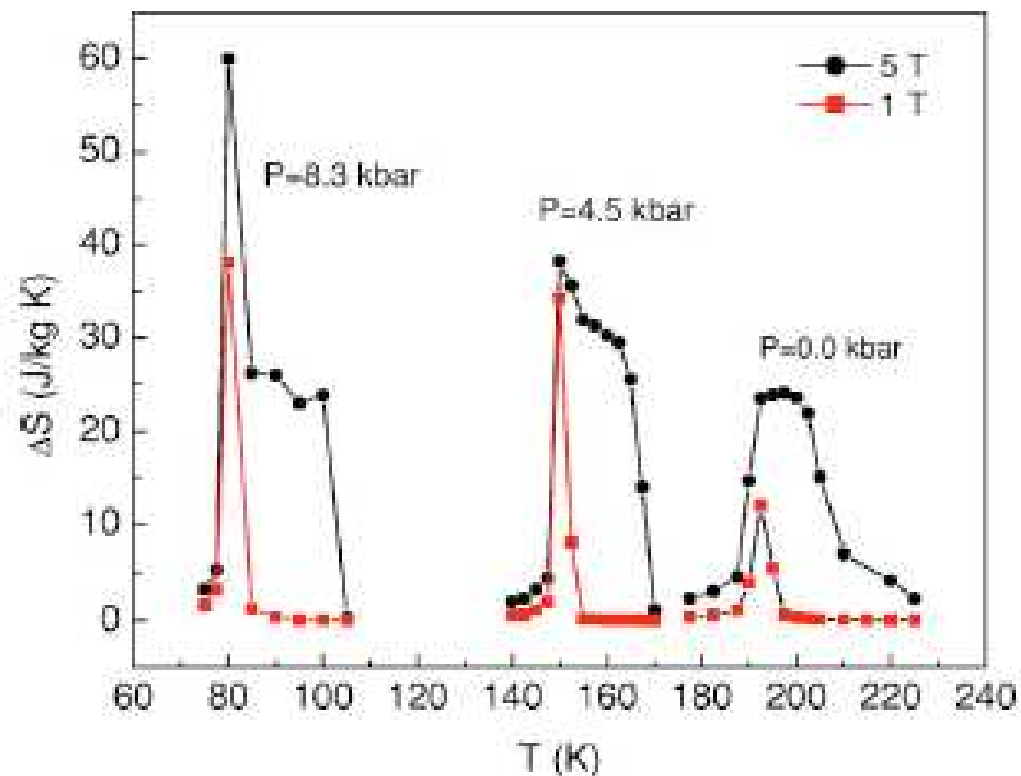


$$dT_c/dp = 2.64 \text{ K/kbar}$$

Morellon *et al.* PRL 93(2004)



$$dT_c/dp = +3.0 \text{ K/kbar}$$



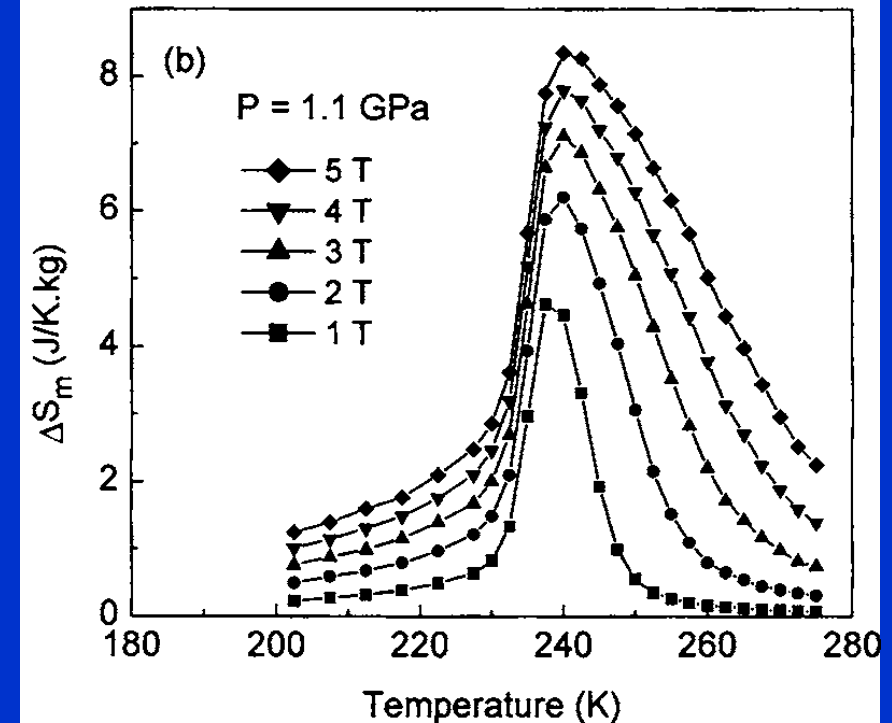
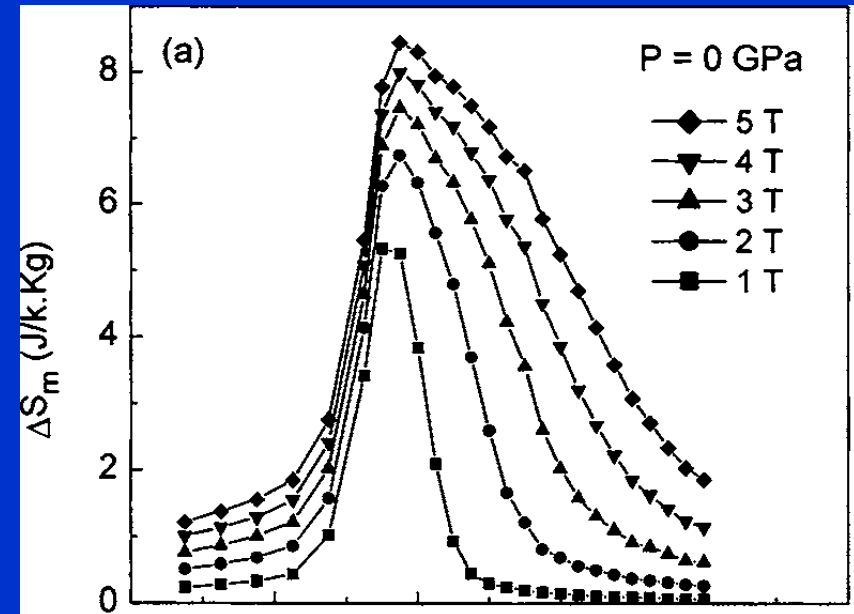
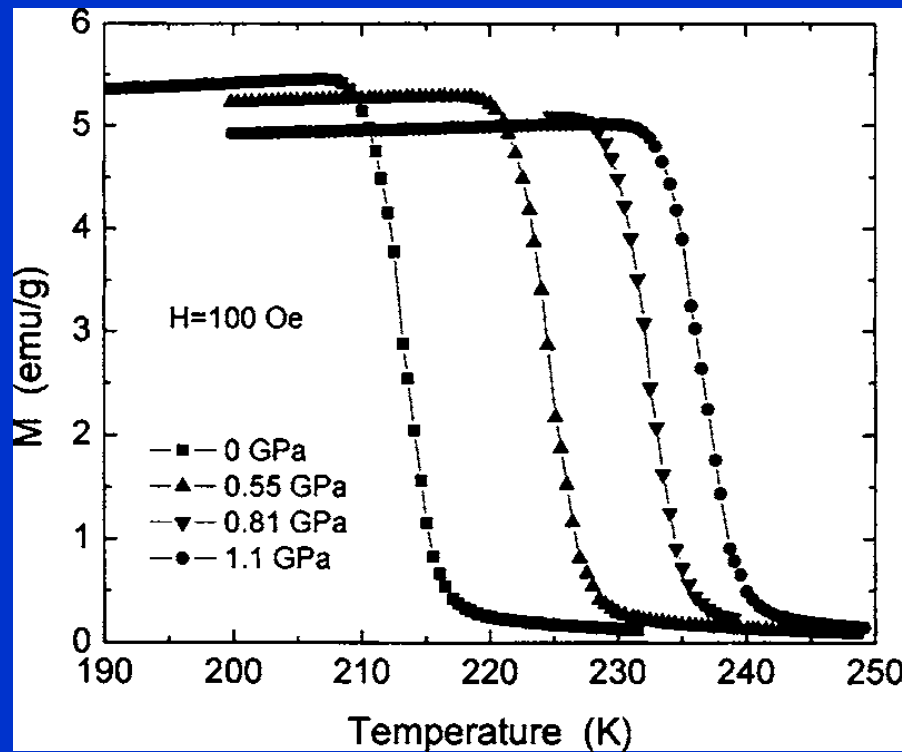
Curie temperature becomes lower!

MCE becomes higher!

(La,Ca)MnO₃

T_h increases with p

Y. Sun et al. APL 88 (2006)



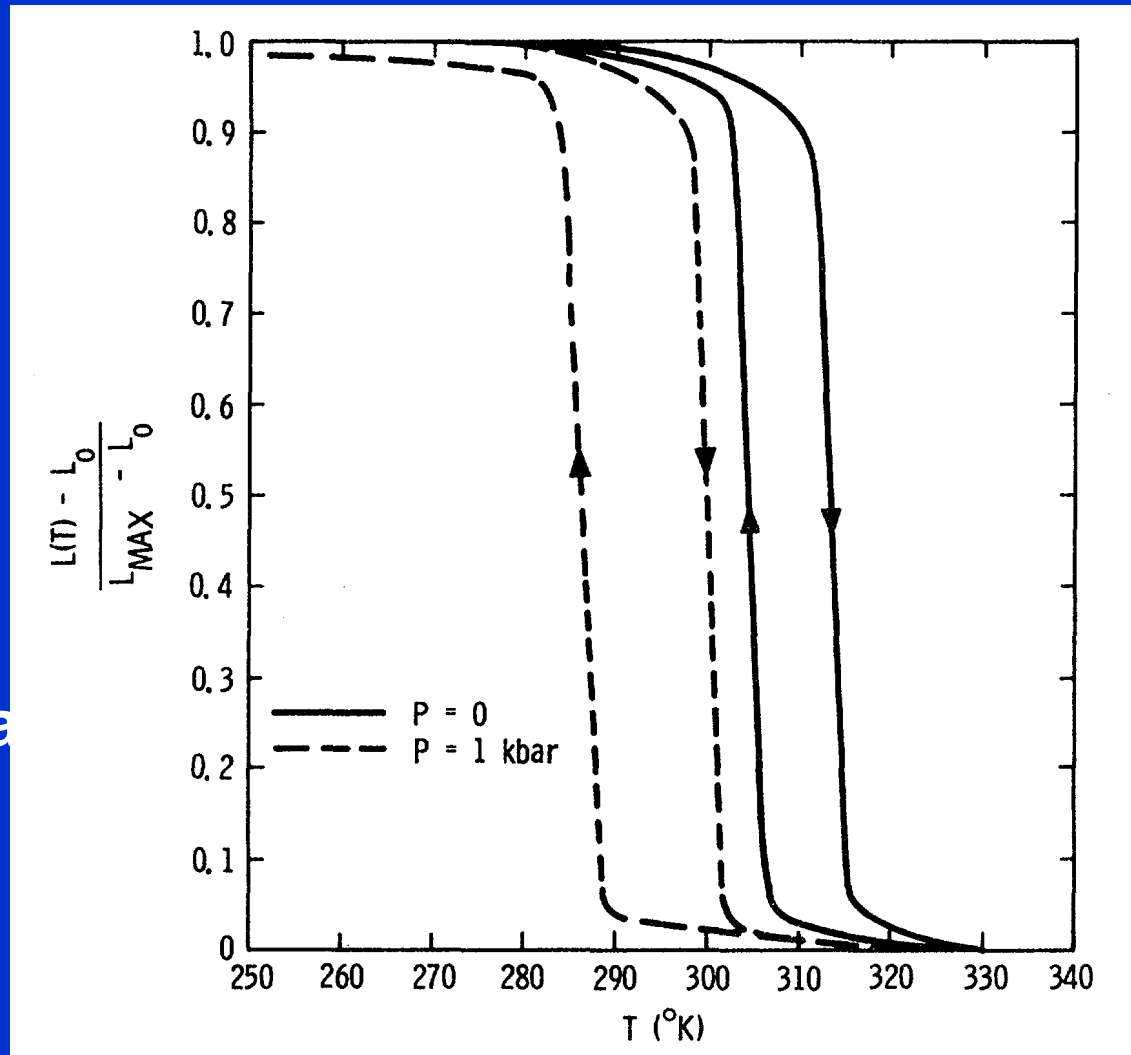
$\text{MnAs}_{0.9}\text{Sb}_{0.1}$

Hex. FM



Orthor. PM

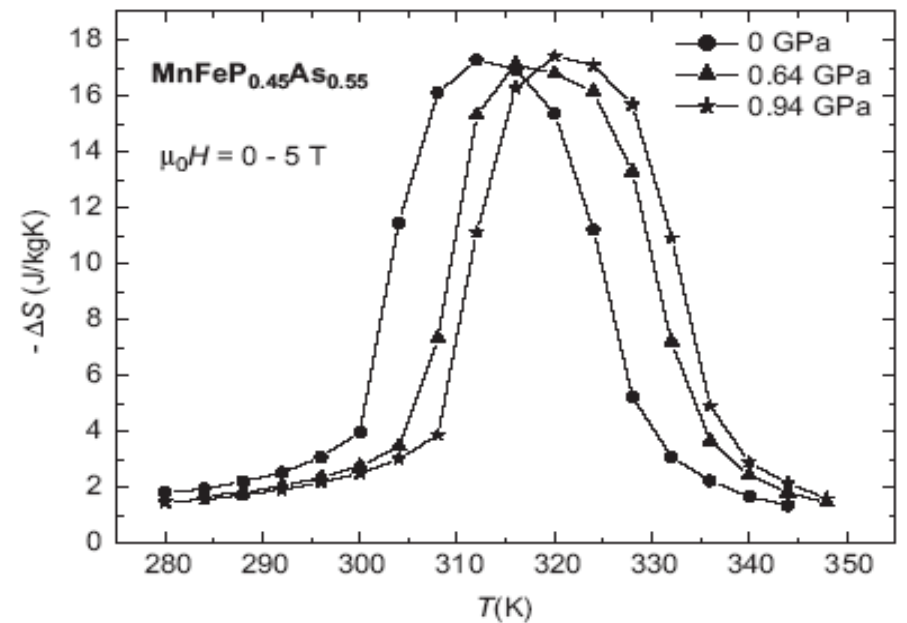
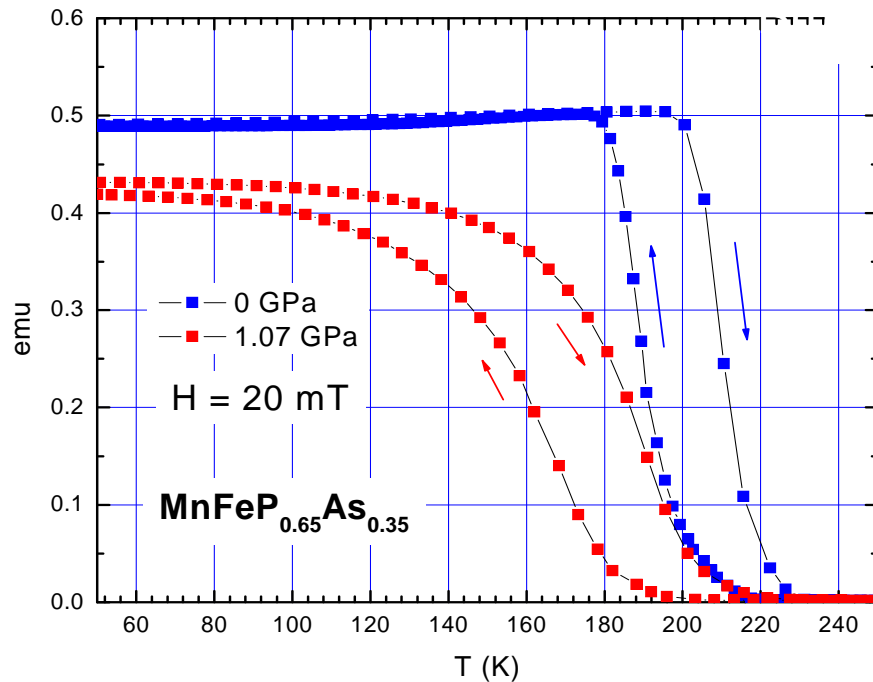
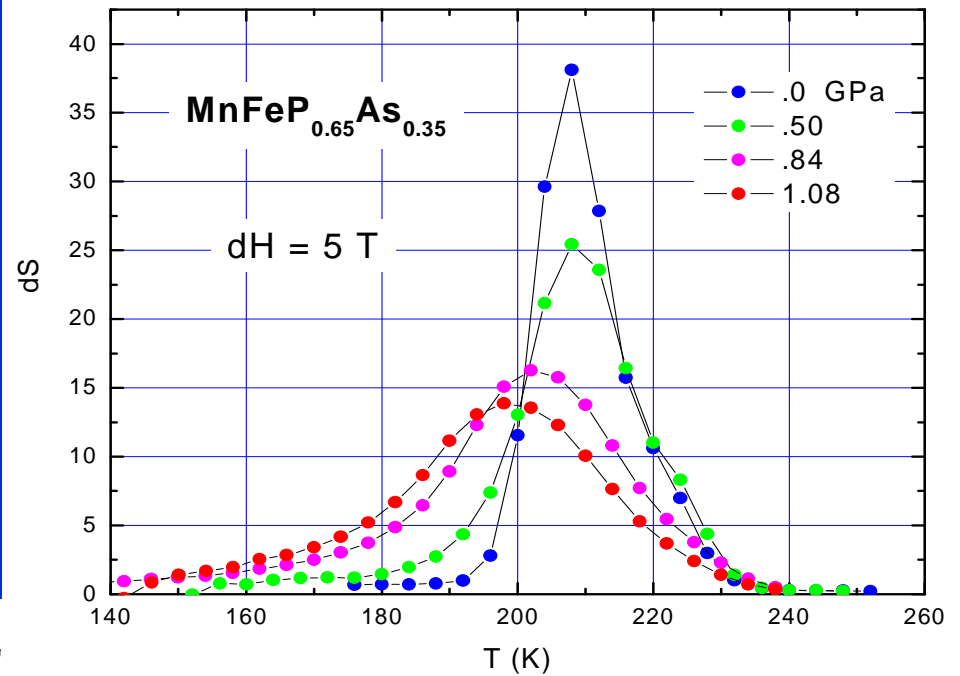
$dT_c/dp = -3.0 \text{ K/kbar}$



Edwards PRB (1972)

Pressure effect on MnFe(P,As) Brück et. al.

MnFeP_{0.65}As_{0.35}
dT_c/dp decreases with p
MnFeP_{0.45}As_{0.55} increases
with p



To summarize

compound	T_c (K)	dT_c/dp (K/GPa)
$Ni_{2.15}Mn_{0.85}Ga$	335	4.7
$ErCo_2$	33	-16.6
$Gd_5Si_2Ge_2$	276	30
$MnAs_{0.9}Sb_{0.1}$	>300	-30
$La_{0.69}Ca_{0.31}MnO_3$		21.4
$La(FeSi)_{13}$		-90
$MnFe(P,Ge)$		negative

dT_c/dp depends on pressure

Thermal hysteresis also depends on pressure.

MnFeP_{1-x}As_x compounds

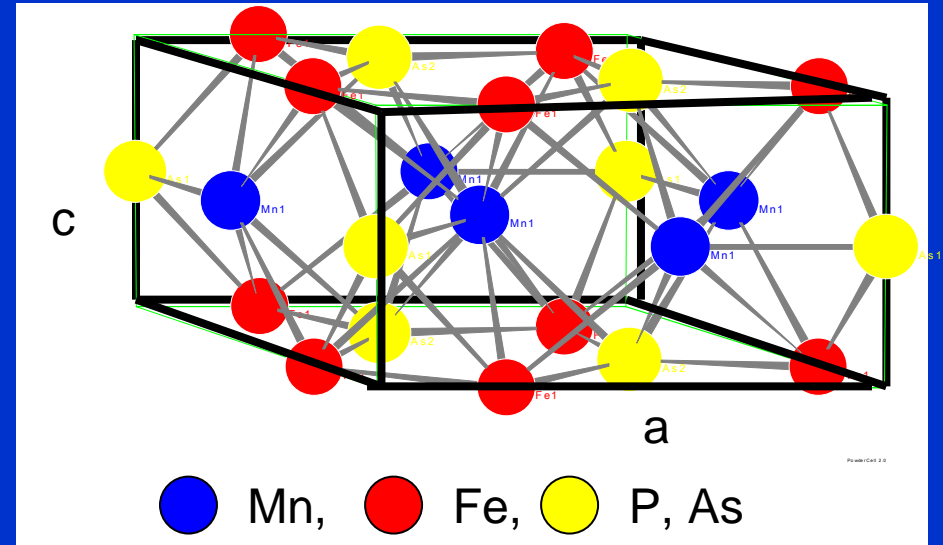
■ MnFeP_{1-x}As_x (0.2 < x < 0.66)

- Hexagonal Fe₂P-type structure

Zach et al. (1990)

- A first order transition from FM to PM state with a hysteresis at T_C.

- No structural change



■ X ~ 0.5: T_C ~ RT

■ A large magnetocaloric effect at RT --> magnetic refrigerant

Tegus et al. Nature 415 (2002) 150.

MnFeP_{1-x}As_x compounds

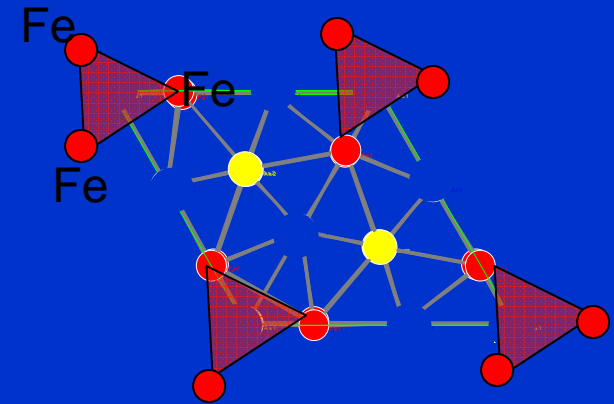
Magnetic moments in MnFeP_{0.5}As_{0.5}

Neutron diffraction Bacmann et al. (1994)

Magnetic moments at 200 K

$$\text{Mn: } m_{\text{Mn}} = 2.02 \mu_{\text{B}}$$

$$\text{Fe: } m_{\text{Fe}} = 1.48 \mu_{\text{B}}$$



in c-plane

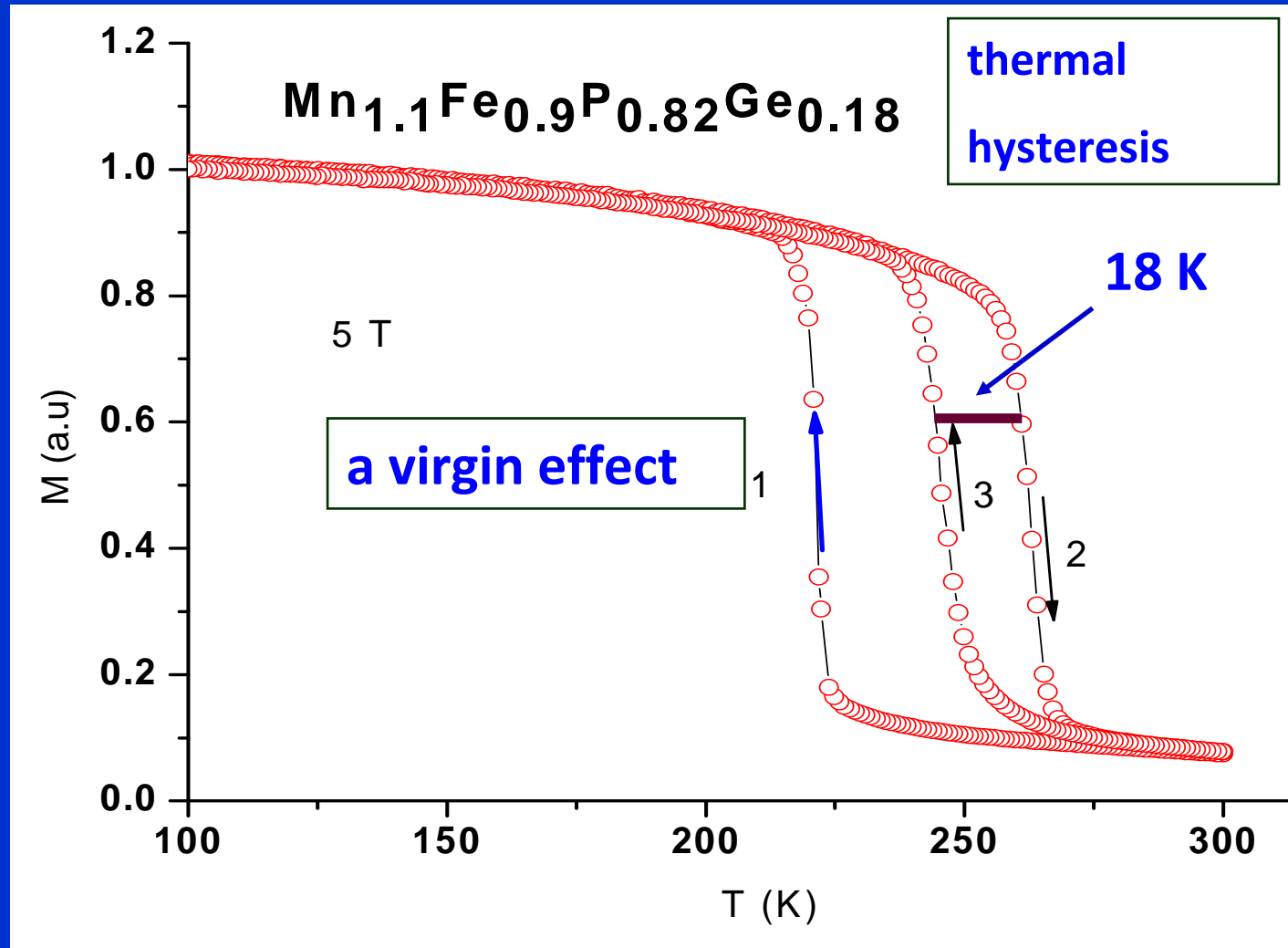
Band structure calc. by Yamada & Terao (2002)

- hybridization of Fe(3d)-Fe(3d) in c-plane
--> suppresses (unstable) m_{Fe}
- (interlayer) hybridization of Fe(3d)-Mn(3d)
--> enhances (stable) m_{Fe}

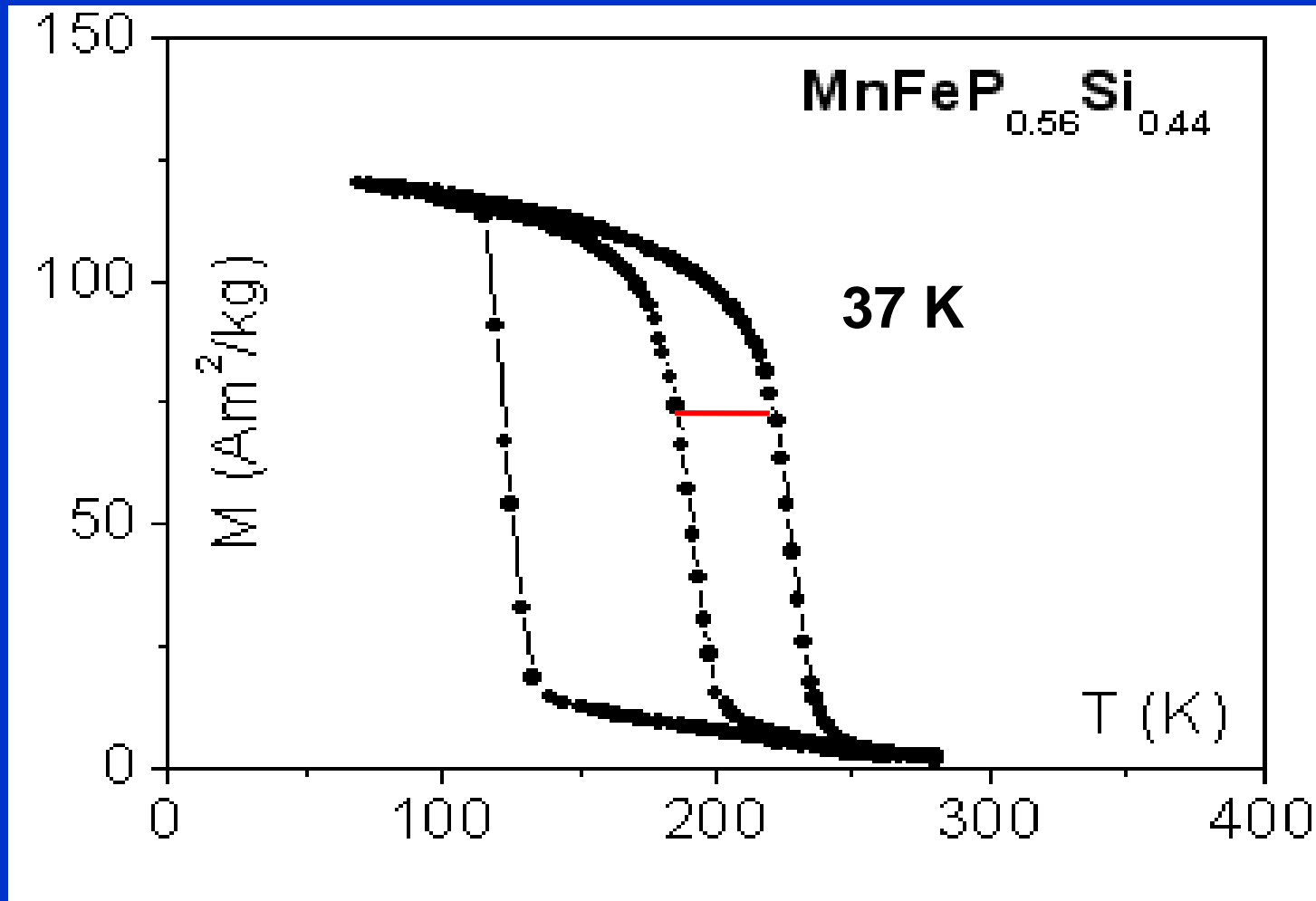
Applying pressure

a decreases and c increases \rightarrow FM state becomes unstable

Problems & Purpose of this study



Cam
Thanh
et. al.



1. New phenomenon: low transition point, irreversible, why?
2. Large thermal hysteresis, what is the mechanism?

Theoretical model

including the coupling between the strain and magnetization

Gibbs free energy of the system

$$g = g_0 + A\sigma^2 + B\sigma^4 + C\sigma^6 + D\omega\sigma^2 + E\omega\sigma^4 + p\omega + \frac{1}{2K}\omega^2$$

σ : spontaneous magnetization,

ω : relative volume change,

K : the compressibility,

p : applied pressure.

$$\omega = \frac{V - V_0}{V_0}$$

A , B , and C are the Landau coefficients

D and E represent coefficients of the third and fifth order coupling between strain and magnetization.

The equilibrium conditions

$$\frac{\partial g}{\partial \sigma} = 0$$

$$A + 2(B + E\omega)\sigma^2 + 3C\sigma^4 + D\omega = 0$$

$$\frac{\partial g}{\partial \omega} = 0$$

$$p + \frac{\omega}{K} + D\sigma^2 + E\sigma^4 = 0.$$

The solution of σ

$$\sigma^2 = \frac{-B' + (B'^2 - 4A'C')^{1/2}}{2C'}$$

where

$$A' = A - DKp$$

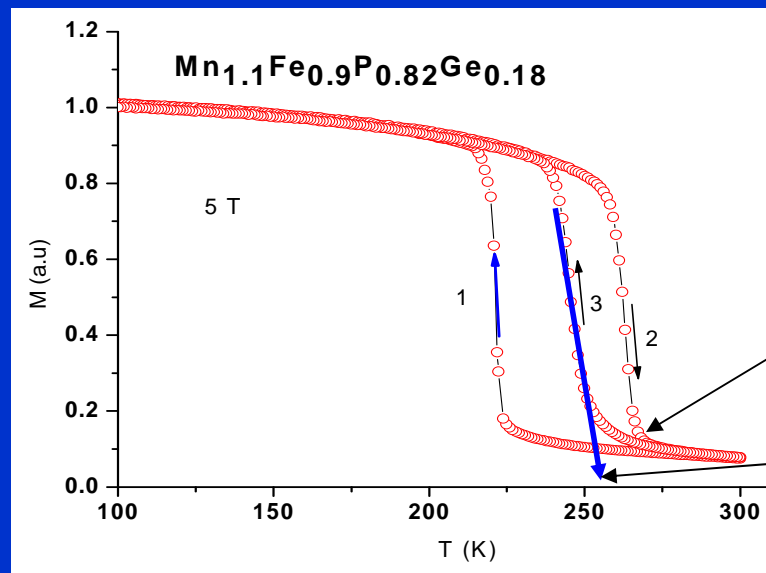
$$B' = 2B - D^2K - 2EKp$$

$$C' = 3(C - EDK)$$

T_0 - At which the theoretical magnetic susceptibility becomes infinite;

T_h - the highest temperature for which the ferromagnetic phase can exist as a metastable state

T_c - the lowest temperature for which the paramagnetic phase can exist as a metastable state.



T_h

T_c

The high-temperature paramagnetic phase remains metastable until $A' = 0$, taken $A = A_0(T-T_0)$

$$T_c = T_0 + \frac{DK}{A_0} p$$

$$\frac{dT_c}{dp} = \frac{DK}{A_0}.$$

T_c linearly depends on pressure

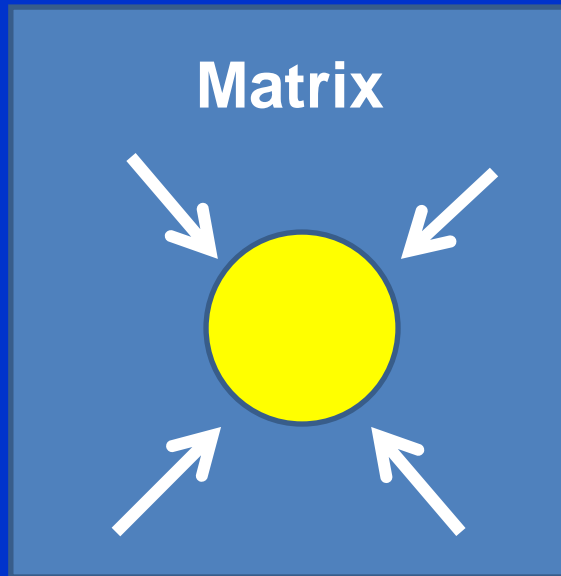
If $DK > 0$ T_c increases with p

If $DK < 0$ T_c decreases with p

If $D = 0$ $T_c = T_0$, no p dependent.

Experiment show that T_c depends on P , therefore, $D \neq 0$, indicating third-order coupling between strain and magnetization should exist.

A possible explanation of the virgin effect



$$T_c = T_0 + \frac{DK}{A_0} p$$

$$DK < 0$$

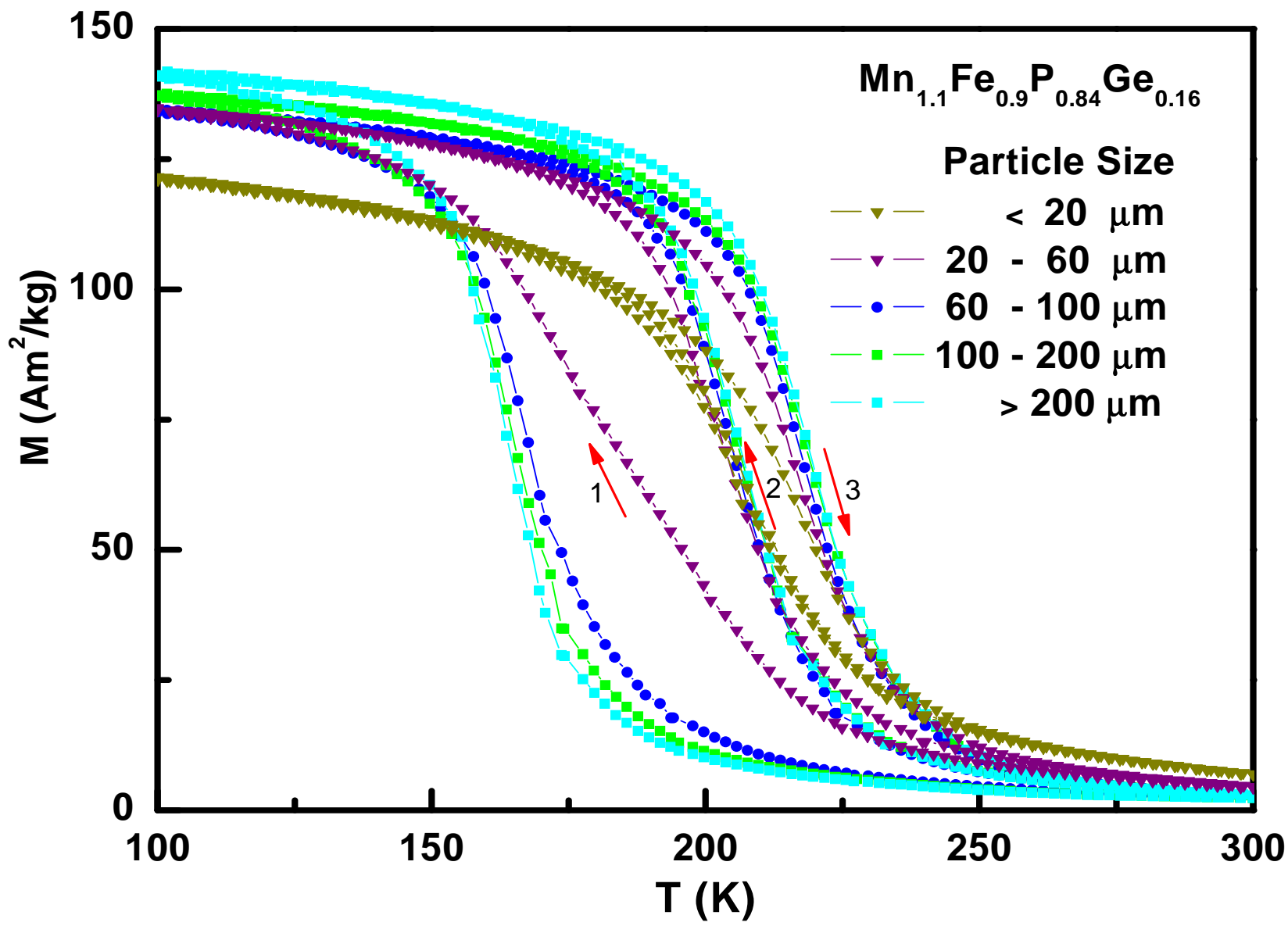
Misfit strain exists.

Additional elastic energy

Product phase becomes powder, or occurs microcracks

Release misfit strain, disconnected

Only exist in cooling, and once!



The low-temperature ferromagnetic phase remains metastable until

$$B'^2 = 4A'C'$$

$$T_h = T_c + \frac{(2B - D^2K - 2EKp)^2}{12A_0(C - EDK)}$$

T_h is not linearly dependent on p .

If $E = 0$

$$T_h = T_c + \frac{(2B - D^2K)^2}{12A_0(C - EDK)} = T_0 + \frac{(2B - D^2K)^2}{12A_0C} + \frac{DK}{A_0} p$$

Then T_h also linearly depends on p

The spread of the thermal hysteresis

$$\Delta T = T_h - T_c = \frac{(2B - D^2K - 2EKp)^2}{12A_0(C - EDK)}$$

***P* dependent**

If $E = 0$

$$\Delta T = T_h - T_c = \frac{(2B - D^2K)^2}{12A_0C}$$

***P* independent**

$E \neq 0$; indicates the fifth order coupling between the strain and magnetization should exist.

The pressure dependence of the thermal hysteresis is caused by the fifth order coupling between the strain and the magnetization.

When $p = 0$, then

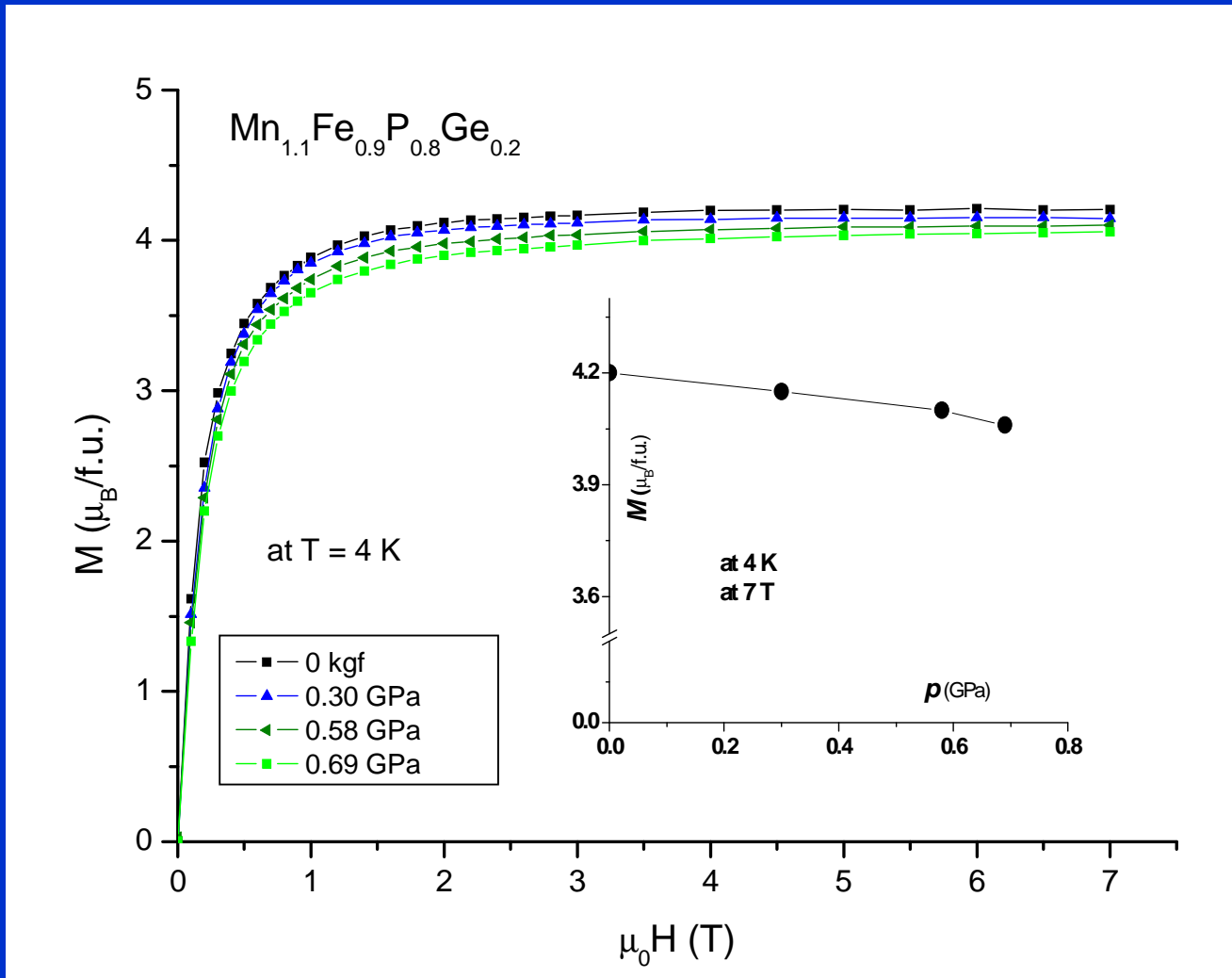
$$\Delta T = T_h - T_c = \frac{(2B - D^2 K)^2}{12A_0(C - EDK)}$$

Without coupling between strain and magnetization

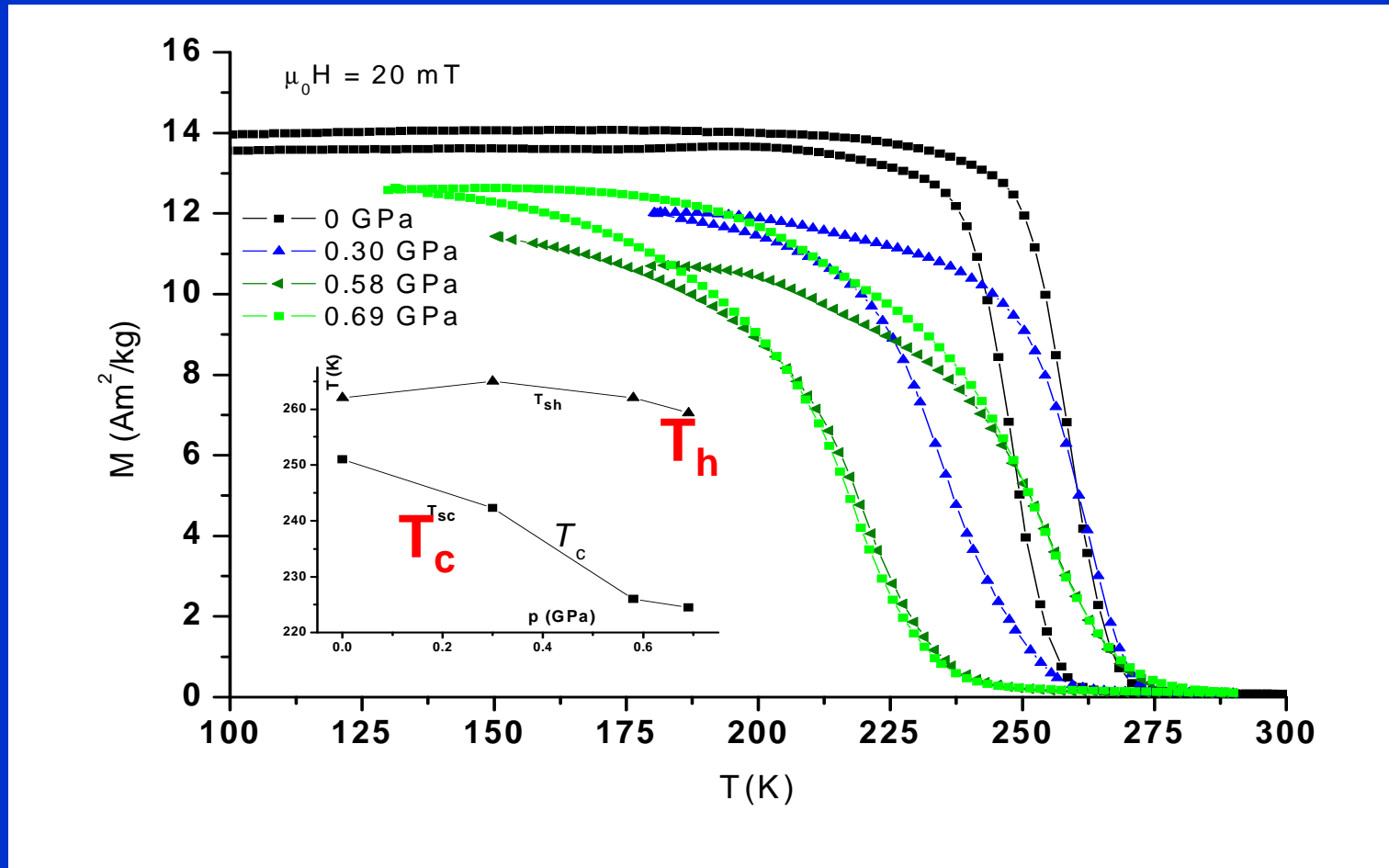
$$\Delta T = T_h - T_c = \frac{B^2}{3A_0 C}$$

Experimental results

Pressure effect on magnetization



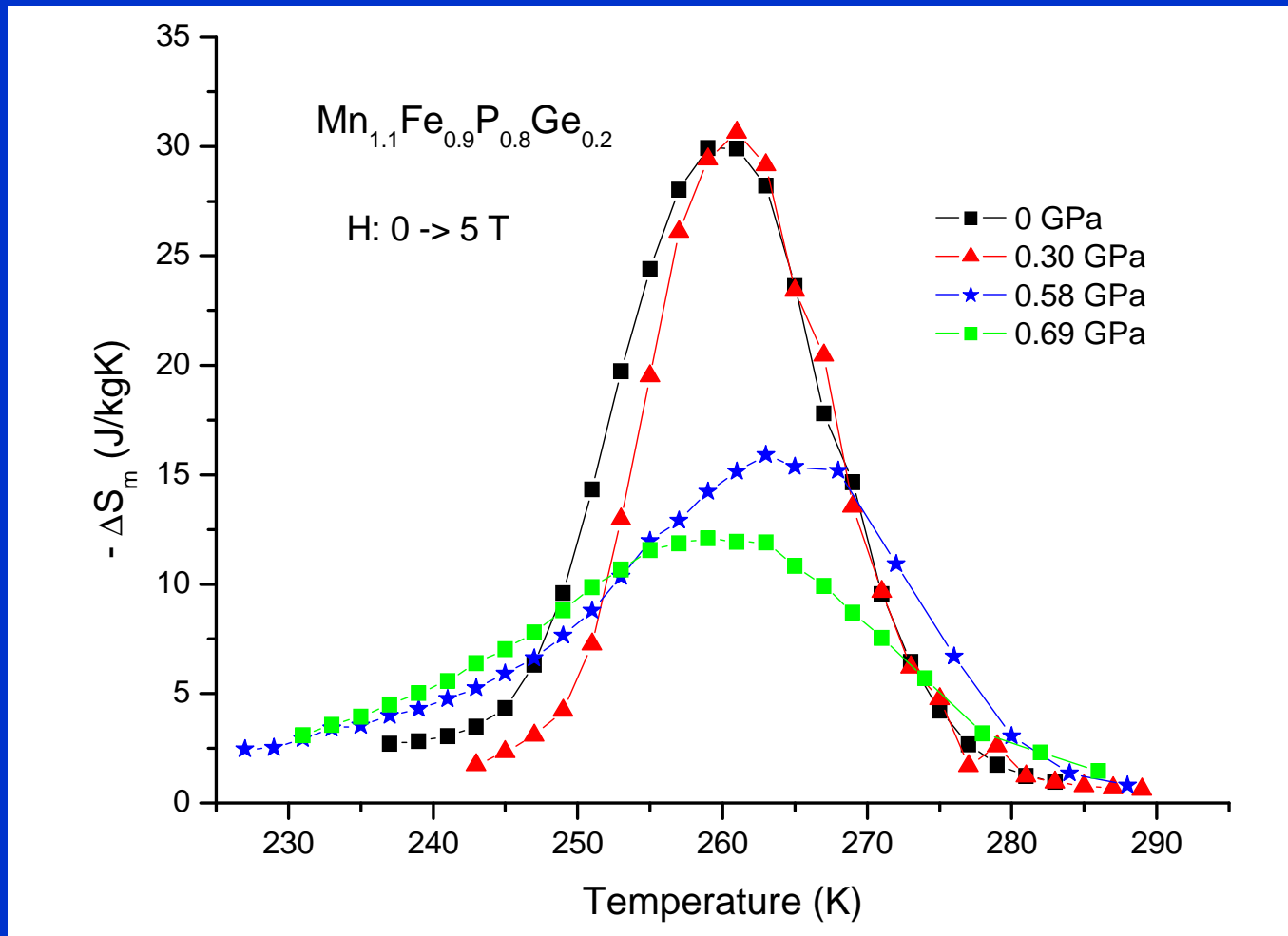
Pressure effects on thermal hysteresis



$$T_c = T_0 + \frac{DK}{A_0} p$$

$$T_h = T_0' + \alpha p + \beta p^2$$

Pressure effects on MCE



Conclusions

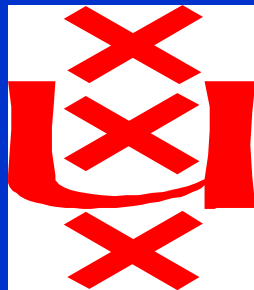
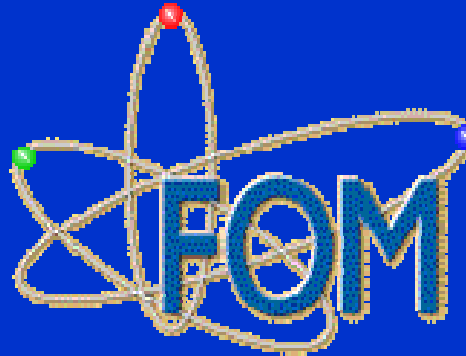
1. The T_c , M , ΔS_m decrease with increasing pressure
2. The thermal hysteresis increases.
3. These pressure dependence of the behaviors can be understood in the framework of Landau theory, but the Gibbs free energy should include the third and fifth order couplings between strain and magnetization.
4. The third order coupling leads to a decrease of T_c and the fifth order coupling results in the increase of thermal hysteresis.

Acknowledgments

E. Brück ,
F.R. de Boer,
K.H.J. Buschow
A. Magnus
B. G. Carvalho,
C.A. A. Coelho,
D.S. Gama
Cam Thanh
Z.Q. Ou

KNAW

China Natural Science Foundation



Thank you for your attention!